Plasmonic Enhancement of Graphene Long-Wave Infrared Photodetectors via Bull's Eye Concentrator, Optical Cavity and Nanoantennas

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The long-wave infrared photodetector device we study in this work is shown in Fig. 1 (a) and (b). A bull's eye grating collects the incident light from a large area ($101 \mu\text{m}$), and couples it into surface plasmon polaritons (SPPs) [1] that converge at its center, at which there is an opening in the Au plate on which the grating sits. Within this opening (or aperture), Au nanobar antennas separated by small gaps are arranged in a radial fashion. These enhance the electromagnetic field in the near-field zone, particularly in the gaps, through localized surface plasmon resonances (LSPRs) [2]. This enhances the absorption in the underlying material, here a small circle ($6.32 \mu\text{m}$ diam.) of graphene. The absorption is boosted further by forming an optical cavity under the graphene. We optimized the geometrical parameters to maximise absorption in the graphene at $\lambda = 9.4 \mu\text{m}$ via finite-difference time-domain electromagnetic simulations. These geometric parameters are listed in Fig. 1 (c)-(f).

Fig. 1 (a) Cartoon illustrating the proposed LWIR photodetector consisting of a hybrid plasmonic structure. (b) Zoomed view of the aperture nanobar antennas. (c)-(f) Geometrical parameters of the structure.

Fig. 2 (a) presents the electric field enhancement at points E_1 and E_2 , whose locations are defined in Fig. 1 (d). We observe an electric field enhancement of ~ 380 and ~ 150 times at E_1 and E_2 , respectively, at the resonance wavelength ($\lambda = 9.4 \mu\text{m}$). Fig. 2 (b) plots normalized power absorbed in graphene vs wavelength, which peaks at $\sim 1.63\%$ at the resonance wavelength. In this figure, the normalized power absorbed is the power absorbed in the graphene divided by the power incident through a square window with a side length of $101 \mu\text{m}$.

Fig. 2 (a) Electric field enhancement at points E_1 and E_2 , respectively. E_i is the norm of the incident electric field. (b) Normalized power absorbed in graphene as a function of wavelength. (c) and (d) Schematic of the non-plasmonic device on a conventional SiO_2/Si substrate with a small-area and large-area graphene flake, respectively.

Our structure is predicted to enhance the absorption of light by the graphene by ~ 4600 times compared with its non-plasmonic counterpart (Fig. 2 (c)). It is predicted to enhance detectivity of the LWIR photodetector by ~ 288 times compared with its non-plasmonic graphene counterpart (Fig. 2 (d)). This estimate is made using $D^* = \sqrt{A_D \Delta f} / NEP$ [3]. This assumes that the photoactive material is coupled to the incident light by an optical concentrator with area A_O , with Δf and NEP being bandwidth and noise equivalent power. NEP is assumed to scale as $\sqrt{A_D} / (Gr \text{ abs})$, where A_D is graphene area and $Gr \text{ abs}$ is normalized power absorbed in the graphene.

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