X-RAY PROPERTIES OF THE FIRST SZE-SELECTED GALAXY CLUSTER SAMPLE FROM THE SOUTH POLE TELESCOPE


ABSTRACT

We present results of X-ray observations of a sample of 15 clusters selected via their imprint on the cosmic microwave background (CMB) from the thermal Sunyaev-Zel’dovich (SZ) effect. These clusters are a subset of the first SZ-selected cluster catalog, obtained from observations of 178 deg$^2$ of sky surveyed by the South Pole Telescope. Using X-ray observations with Chandra and XMM-Newton, we estimate the temperature, $T_X$, and mass, $M_X$, of the intracluster medium (ICM) within $r_{500}$ for each cluster. From these, we calculate $Y_X = M_X / T_X$ and estimate the total cluster mass using a $Y_{500} - Y_X$ scaling relation measured from previous X-ray studies. The integrated Comptonization, $Y_{SZ}$, is derived from the SZ measurements, using additional information from the X-ray measured gas density profiles and a universal temperature profile. We calculate scaling relations between the X-ray and SZ observables, and find results generally consistent with other measurements and the expectations from simple self-similar behavior. Specifically, we fit a $Y_{SZ} - Y_X$ relation and find a normalization of $0.82 \pm 0.07$, marginally consistent with the predicted ratio of $Y_{SZ}/Y_X = 0.91 \pm 0.01$ that would be expected from the density and temperature models used in this work. Using the $Y_X$ derived mass estimates, we fit a $Y_{SZ} - M_{500}$ relation and find a slope consistent with the self-similar expectation of $Y_{SZ} \propto M^{5/3}$, with a normalization consistent with predictions from other X-ray studies. We compare the X-ray mass estimates to previously published SZ mass estimates derived from cosmological simulations of the SPT survey. We find that the SZ mass estimates are lower by a factor of $0.89 \pm 0.06$, which is within the $\sim 15\%$ systematic uncertainty quoted for the simulation-based SZ masses. Overall, the X-ray measurements confirm that the SZ-selected sample consists of very massive systems which exhibit general properties consistent with X-ray selected samples, even allowing for the broad redshift range ($0.29 < z < 1.08$) of the sample.

Subject headings: galaxies: clusters — X-rays: galaxies: clusters

1. INTRODUCTION

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Large area cluster surveys extending to high redshift can be used to study the evolution of the abundance of galaxy clusters, thereby delivering precise constraints on the amount and nature of the dark energy (Wang & Steinhardt 1998; Haiman et al. 2001). The accuracy with which the observed mass proxy can be linked to the true cluster mass fundamentally limits the cosmological constraints from the survey. In particular, a redshift dependent bias on a survey’s cluster mass estimates could mimic a time-evolving dark energy, so this systematic must be understood and constrained. Clusters are known to evolve — through mergers, galaxy and star formation, and variable contributions from active galactic nuclei (AGN) — and these effects will influence the evolution of any cluster observable at some level.

Most of the baryonic mass in clusters is in the form of an intra-cluster gas that can be heated to several keV as it virializes in a massive cluster’s gravitational potential well. This gas is visible in the X-ray band via thermal bremsstrahlung and also from its distortion of the Cosmic Microwave Background (CMB) from inverse Compton scattering, otherwise known as the Sunyaev-Zel’dovich (SZ) effect (Sunyaev & Zeldovich 1972). The lowest scatter cluster observables that scale with cluster mass, $M$, are likely those most closely related to the gas pressure (e.g., Kravtsov et al. 2006), and hence related to the total energy of the gas. The SZ intensity is proportional to the Comptonization, the line-of-sight integral of the gas pressure. Hence, the SZ signal integrated over the cluster’s extent, $Y_{SZ}$, measures the total pressure in the cluster. An X-ray analogue, $Y_X$, can be constructed from the product of the cluster’s total gas mass and the X-ray spectroscopic temperature (Kravtsov et al. 2006).

The intrinsic scatter in the $M-Y_X$ relation has been found to be smaller than the measurement error in X-ray studies conducted to date (Vikhlinin et al. 2006, 2009a; Sun et al. 2009), and has been used to place interesting constraints on the dark energy equation of state with only a relatively small sample of clusters (Vikhlinin et al. 2009b). The slope and scatter of the relationship between $Y_{SZ}$ and the cluster mass, $M$, has been studied extensively in simulations and is expected to be relatively insensitive to non-gravitational physics (Nagai 2006), the dynamical state of clusters (Jeltema et al. 2008), and the presence of cool cores (Motl et al. 2005). The expected close correlation between cluster mass and $Y_{SZ}$, as well as the redshift independence of the SZ brightness, strongly motivates using SZ cluster surveys for cosmological studies (Carlstrom et al. 2002).

Recently, there has been significant progress in measuring the SZ-signal from clusters. High resolution imaging has been obtained for single objects (e.g., Nord et al. 2009; Mason et al. 2010) and ICM profiles have been measured for moderately sized samples (e.g., Mroczkowski et al. 2009; Plagge et al. 2010). The first clusters discovered with a blind SZ survey were reported in Staniszewski et al. (2009), and showed the capability of the SZ-signal as a cluster finder. From larger samples, scaling relations have been measured between the SZ-signal and mass estimates from both X-ray and weak lensing measurements (e.g., Bonamente et al. 2008; Marrone et al. 2009; Melin et al. 2010). Notably, Melin et al. (2010) measured a $Y_{SZ}-M$ relation from binned WMAP fluxes at the location of known X-ray selected clusters combined with X-ray luminosity based mass estimates. They found a $Y_{SZ}-M$ relation with a normalization and slope that matched the X-ray prediction, however similar analyses have found conflicting results (Komatsu et al. 2010). More detailed observations comparing SZ and X-ray measurements will improve our understanding of the gas pressure in clusters.

The first SZ-selected catalog of clusters was presented in Vanderlinde et al. (2010) (hereafter V10), and included the first meaningful cosmological constraints from an SZ cluster survey. In V10, the sample of 21 clusters had a median redshift of $z = 0.74$ and was predicted to be 100% complete above a mass threshold of $M_{500} \approx 3 \times 10^{14} h^{-1} M_{\odot}$ at $z > 0.6$. The cosmological constraints from V10 were limited by uncertainties in the cluster mass calibration. This calibration relied on the dark matter simulations of Shaw et al. (2009), with gas physics based on the models in Bode et al. (2007), to link the SZ-significance to cluster mass. This introduced a ~15% systematic uncertainty in the mass calibration due to uncertainties in the pressure normalization of the simulations. Therefore, an important first step to improve the cosmological constraints of V10 is to tie the SZ observables to observationally calibrated X-ray scaling relations, such as those in Vikhlinin et al. (2009a).

The V10 cluster sample is also unique because of its high median redshift and SZ-selection. X-ray observations of these clusters could allow additional constraints on the redshift evolution of X-ray scaling relations, which typically have been studied from X-ray selected samples concentrated at $z < 0.6$ (e.g., Maughan 2007; Vikhlinin et al. 2009a; Mantz et al. 2009).

In this work, we present results from X-ray observations of a subset of 15 clusters with the highest SZ-significance from V10. We report on the X-ray observables of this sample, and use X-ray scaling relations from Vikhlinin et al. (2009a) to estimate each cluster’s mass. For each cluster, we also note details of the cluster’s X-ray morphology and the identification of X-ray sources with objects in other catalogs. We use the X-ray measured gas density profiles to improve the SZ estimates of integrated Comptonization, $Y_{SZ}$. Finally we construct X-ray and SZ scaling relations, specifically the $Y_{SZ}-Y_X$ and $Y_{SZ}-M_{500}$ relations, and compare these relations to expectations and other results.

The paper is organized as follows: §2 describes the cluster sample and the SZ, optical and X-ray observations. §3 discusses the X-ray data analysis, the estimation of gas mass and temperature while §4 describes the estimate of gravitational mass. In §5 we discuss the analysis of the SPT data and deprojection of the SZ measurements. In §6 the X-ray and SZ scaling relations are investigated and we conclude with a discussion in §7. The properties of individual clusters are discussed in the Appendix.

In all calculations we have assumed a WMAP7+BAO+$H_0$ ΛCDM cosmology (Komatsu et al. 2010) with $\Omega_M = 0.272$, $\Omega_{\Lambda} = 0.728$ and $H_0 = 70.2$ km s$^{-1}$ Mpc$^{-1}$ with distance measurements from Baryon Acoustic Oscillations (BAO) in the distribution of galaxies (Percival et al. 2010) and the Hubble constant ($H_0$) measurement from Riess et al. (2009). Everywhere, we define $M_{500}$ as the mass inside $r_{500}$, within which the matter density is 500 times the critical density at the cluster redshift, $\rho_{\text{crit}}(z) = 3H^2(z)/8\pi G$, where $H(z)$ is the Hubble parameter.
2. CLUSTER SAMPLE AND OBSERVATIONS

The clusters used in this work were selected from a SZ detection significance-limited catalog from the South Pole Telescope (SPT) cluster survey described in V10. For all clusters, we performed follow-up optical imaging to identify galaxy cluster counterparts and measure redshifts (High et al. 2010). Sixteen cluster candidates with SZ detection significance above 5.4 were selected for X-ray follow-up with Chandra and XMM-Newton. One candidate was later discovered to be a spurious detection, and is discussed in more detail in §A. This section briefly describes the SZ, optical, and X-ray data sets.

2.1. SPT SZ Observations and the Cluster Sample

The 10-m diameter South Pole Telescope is a millimeter-wavelength telescope located at the South Pole. Its primary science goal is to conduct a ∼2000 deg$^2$ survey to find clusters of galaxies via measurements of the SZ effect. The receiver consists of a 960 element bolometer array that is sensitive in three frequency bands centered at 95, 150, and 220 GHz. In this work, we only use observations at 150 GHz, the SPT frequency band with the most SZ sensitivity. Details of the telescope and receiver can be found in Padin et al. (2008) and Carlstrom et al. (2009).

In 2008, the SPT surveyed two ∼100 deg$^2$ regions. These two fields are approximately square on the sky, and centered at right ascension (R.A.) 5$^h$30$^m$, declination (decl.) -55$^\circ$ and R.A. 23$^h$30$^m$, decl. -55$^\circ$. The data from the first of these fields was used to report the first SZ-discovered clusters (Staniszewski et al. 2009), to measure source counts of extragalactic mm-wavelength emitting objects (Vieira et al. 2009), and to measure small scale temperature anisotropies due to the SZ effect from unresolved clusters and emission from point sources (Lueker et al. 2009; Hall et al. 2009). In V10, the data from both fields was used to report the first SZ detection significance-limited catalog of 22 cluster candidates.

The SZ observations, data processing, and mapmaking used in this work are described in detail in V10, and only a brief overview is provided here. Each field was scanned in azimuth at a constant velocity, with the scans stepped in elevation. With the velocity and elevation step used, it takes ∼2 hours to cover an entire 100 deg$^2$ field. This scan is then repeated several hundred times to decrease the noise in the co-added map. Each detector timestep is filtered to remove both long timescale drifts and sky signal that is spatially correlated across the focal plane. The filtering effectively acts as a high pass filter; the filter signal that is spatially correlated across the focal plane.

Cluster candidates were identified in the SPT maps by using a matched spatial filter technique (Haehnelt & Tegmark 1996; Melin et al. 2006). Here we summarize the method and results as used in V10. The SPT maps consist of several sources of signal, including primary CMB anisotropy, unresolved point sources, and SZ signal from clusters. To identify cluster candidates, the SPT maps are filtered in Fourier space to give more weighting to signals matching the expected spatial scales of clusters. Twelve different spatial filters were constructed using spherical $\beta$-models with $\beta$ fixed to 1 and core radii evenly spaced between 0.25$'$ and 3.0$'$. With these, the SPT maps were filtered and cluster candidates were identified as decrements in the filtered map. The significance of a candidate was quantified by their signal relative to the standard deviation in the filtered map, or signal-to-noise. For a given cluster candidate, the highest signal-to-noise across all filter scales was defined as $\xi$. To avoid spurious identifications from ringing around spatially filtered bright point sources, a 4$'$ radius region around all 5$\sigma$ positive sources was masked before spatial filtering. The total sky area used in both fields after masking was 178 deg$^2$. Further details of the SZ analysis relevant to the results presented here are given in §5.

The sixteen cluster candidates with the largest $\xi$ were selected for X-ray follow-up with the Chandra and XMM-Newton X-ray telescopes. The original candidate list was based on a similar, but earlier version of the list that appeared in V10. This change re-ordered the list somewhat, such that sixteen of the seventeen most significant candidates from V10 had X-ray follow-up. In Table 1, we give the position and $\xi$ from V10 for the fifteen confirmed clusters with X-ray observations. One of the candidates with X-ray follow-up, SPT-CL J2343-5521, is not listed because it is very likely a false detection. Its X-ray observation is discussed further in the Appendix.

2.2. Optical Imaging and Spectroscopy

Optical counterparts for the clusters selected by SZ detection significance were identified and photometric redshifts were measured via a combination of imaging from the Blanco Cosmology Survey (BCS) and targeted observations using the Magellan telescopes. For a subset of clusters, spectroscopic redshifts were also obtained. Further details of the optical data and analysis can be found in High et al. (2010), and are briefly described below.

The BCS is an optical survey of two ∼50 deg$^2$ fields that lie inside the two 2008 SPT fields described in §2.1. BCS used the Mosaic-II wide field imager on the Blanco 4-m telescope at the Cerro Tololo Inter-American Observatory in Chile. The BCS obtained contiguous deep optical imaging in the griz bands across their survey fields, and these data have been processed using a development version of the DES data management system (Ngew et al. 2006; Mohr et al. 2008) and then used in the study of the galaxy population and in the redshift estimation of the first SPT survey fields (Ngew et al. 2009; Staniszewski et al. 2009; High et al. 2010; Brodwin et al. 2010; Zenteno et al. 2010). These data are publicly available through the NOAO Survey Program and have been used by other groups to study clusters in the SPT survey region (Menanteau et al. 2009; Menanteau & Hughes 2009; Menanteau et al. 2010; McMahon et al. 2009; Suhada et al. 2010). For the nine clusters outside the BCS fields, optical imaging in the griz bands was obtained with the Inamori Magellan Areal Camera and Spectrograph (IMACS; Dressler et al. 2006) on the Magellan Baade 6.5-m telescope in Las Campanas, Chile. Five cluster candidates in the BCS fields were also re-observed with Magellan. In contrast to BCS, the Mag-
2.3. Chandra & XMM-Newton observations

As described in §2.1, the sixteen cluster candidates with the highest detection significance, $\xi$, were selected for an X-ray follow-up program. This program was split between several observing cycles and proposals between the Chandra and XMM-Newton X-ray satellites.

In the original planning of X-ray observations, the observing time required for each cluster was estimated from preliminary predictions of the SZ significance-to-mass relation and a mass-luminosity scaling. The uncertainty in the relations does not allow an accurate estimate of the exposure needed for a required number of photons. The estimates were informed by ROSAT fluxes where available. For clusters scheduled for multiple observations with Chandra, the integration time was modified accordingly once a flux had been measured.

To date, the Chandra program consists of two GTO programs in AO-9 (295 ks, total), two GTO programs in AO-11 (340 ks, total) and a GO program in AO-11 (310 ks). When completed, the Chandra program will have collected at least 1500 cluster photons within $0.5r_{500}$ and in the 0.5-7.0 keV energy band for each of 12 clusters in the sample. This limit was chosen to enable measuring of the ICM temperature, $T_X$, to 15% accuracy. For one very faint object, SPT-CL J0553-5005, the data will not be available. For clusters scheduled for multiple observations with Chandra, the integration time was modified accordingly once a flux had been measured.

Four of the candidates have been observed with XMM-Newton, one of which is an archival observation (SPT-CL J0516-5430, also RXCJ0516.7-5430); the other three are from a 2008 program (SPT-CL J0559-5249, SPT-CL J2332-5358 and SPT-CL J2337-5942) and these are all analyzed in this work. We chose to re-observe these two clusters with Chandra to better identify and remove X-ray point-sources in the analysis. The low redshift and high mass of the clusters made these observations possible with a comparably small amount of observing time.

While we find consistent results from the analysis of data from the two different satellites, we choose to use the Chandra data for the two clusters with data from both instruments since both the high spatial resolution and stable background of Chandra are desirable. For the two clusters where only XMM-Newton data were available, we derive our results from these data using the methods as described below.

X-ray measurements of SPT-CL J2332-5358 and SPT-CL J2342-5411 were also reported by Suhada et al. (2010) from the XMM-BCS survey. These measurements were not as deep as the observations discussed in our work, however we find X-ray observables consistent with their measurements.

Currently, there are eight out of the fifteen clusters for which the required 1500 source counts have not yet been collected with either instrument. However, we include these in the analysis since the existing data provide useful constraints.

3. X-ray Data Analysis

3.1. Data reduction

The exposure times and resulting source counts from the Chandra and XMM-Newton observations are listed in Table 1. We also list the cluster coordinates and SZ signal-to-noise as described in V10 as well as the optical redshifts and the XMM-Newton X-ray measurements of SPT-CL J2332-5358 and SPT-CL J2342-5411 were also reported by Suhada et al. (2010) as described below.

The Chandra data were reduced using CIAO 4.1 and CALDB 4.1.3. All data were taken with the ACIS-I nominal aimpoint in VFAINT telemetry mode and additional screening to reject particle background was applied. To remove periods of flaring background we extracted point-source subtracted lightcurves in the 0.3-12 keV band and filtered these using a 3$\sigma$ threshold. XMM-Newton data were reduced using SAS 9.0 and reprocessed. Source-free light curves were generated in hard (MOS:10-12 keV, pn:12-14 keV) and soft (MOS:0.3-10 keV, pn:0.3-10 keV) energy bands separately and a 3$\sigma$ cut was applied to remove periods of high background.

3.2. Data analysis methods

For the X-ray observables, $T_X$, $M_g$ and $Y_X$ we use scaling relations from Vikhlinin et al. (2009a) iteratively to determine the $r_{500}$ radius where the observable is measured.

The $r_{500}$ radii used for the measurement of $T_X$ were estimated using the $M-T_X$ relation in Vikhlinin et al. (2009a) from samples of local clusters with deep exposures for which hydrostatic masses could be determined:

$$M_{500} = (3.02 \pm 0.11) \times 10^{14} \, M_\odot \, h^{-1} \times \left(\frac{kT}{5 \text{ keV}}\right)^{(1.53\pm0.08)} \ E(z)^{-1}$$
we analogously estimate the aperture radius through the 

\[ r_{500} = \left( \frac{3M_{500}}{4\pi 500 \rho_{crit}(z)} \right)^{1/3} \]  

The radius is then defined as

Similarly, in the estimation of the gas mass, \( M_g \), described below, we use the gas-fraction relation

\[ f_{g,500} = \frac{0.0764 \pm 0.004}{h^{-1.5}} \]

\[ + (0.0225 \pm 0.004) h^{-1.5} \log M_{15} \]  

where \( f_{g,500} \) is the gas mass fraction within \( r_{500} \). \( f_{g,500} = M_g/M_{500} \), and \( M_{15} \) is the total mass, \( M_{500} \), in units of \( 10^{15} h^{-1} M_\odot \). With this relation we explicitly take into account the observed trend of \( f_{g,500} \) with cluster mass (e.g., Vikhlinin et al. 2009a). The gas mass estimation is much more dependent on the aperture radius than the estimate of temperature and \( M_g \) must be estimated iteratively to obtain a self-consistent result. Given a gas mass, we use equation 4 to determine the total mass and, hence, \( r_{500} \).

For the estimation of the \( Y_X = M_g \times T_X \) parameter we analogously estimate the aperture radius through the determination

\[ M_{500} = (5.77 \pm 0.20) \times 10^{14} h^{1/2} M_\odot \]

\[ \times \left( \frac{Y_X}{3 \times 10^{-3} M_\odot \text{keV}} \right)^{0.57 \pm 0.03} E(z)^{-2/3} \]  

\[ 26 \text{ This equation contains a typo in Vikhlinin et al. (2009a), Table 3.} \]
complications related to under-subtracted X-ray background which can be a problem when using in-field background. As a second background subtraction method we also use the latest blank field backgrounds collected by M. Markevitch\textsuperscript{27}, re-project these on the sky to match our cluster datasets and normalize the exposure times to match the count rates in the 9.5-12 keV band. The background spectra can then be subtracted from the same detector region as the source. We compare the temperatures and flux estimates from these two techniques and find good agreement. Since the blank-field background subtraction provides better signal-to-noise, we use this method throughout as this allows for temperature estimates within the \( r_{500} \) aperture for most objects.

For the XMM-Newton observations we use the spectral analysis as described in Andersson & Madejski (2004) and Andersson et al. (2009). For this work we limit the analysis to use in-field background since we are only interested in obtaining one spectrum for each cluster. Spectra for both Chandra and XMM-Newton datasets were extracted inside \( r_{500} \), excluding the central emission within 0.15\( r_{500} \) to avoid the effects of cool cores, known to cause additional scatter in X-ray scaling relations. We model the data using a MEKAL model for the thermal plasma with emission lines and a WABS absorption model. For our sample, the data are not deep enough to measure metal abundances reliably and we fix the abundance to 0.3 solar. The hydrogen equivalent absorbing column, \( n_H \), was fixed at the weighted average from the Leiden-Argentine-Bonn survey (Kalberla et al. 2005). These are listed in Table 1. The best fit temperatures are shown in Table 2.

### 3.4. X-ray imaging

The X-ray surface brightness is extracted in \( n \) concentric annuli defined by

\[
r_i = \left( i \times \frac{r_{\text{max}}}{n} \right)^c \quad i = 1 \cdots n,
\]

where we have used \( r_{\text{max}} = 1.5 \ r_{500} \), \( c = 1.2 \) and \( n = 20 \). The limiting radius \( r_{\text{max}} \) is chosen to be large enough so that confusion between background and cluster flux is negligible and so that the integrated gas density is not overestimated due to projection effects. The values of \( c \) and \( n \) are chosen to balance high central resolution with achieving a similar number of X-ray counts per radial bin.

For every annular region the average exposure is calculated, taking into account bad pixels and chip gaps. Similarly, the average effective mirror area is calculated for each region taking into account the mirror vignetting. The radial model for the X-ray counts includes a component for the cluster X-ray surface brightness, a spatially flat unvignetted component for the particle induced X-ray background and soft X-ray background plus an additional spatially flat vignetted component taking into account under-subtraction of the soft X-ray background flux due to mirror vignetting. This last component is always small.

The cluster surface brightness is taken to be proportional to the integrated emission measure, \( EM = \int n_e n_p dV \), which is a good approximation in the 0.5 - 2 keV band for gas temperatures present in our sample. The inferred density distribution has been shown previously to have little dependence on the temperature when determined from the surface brightness in this energy band (Mohr et al. 1999; Forman et al. 2007). The radial gas density distribution is modeled using a modified \( \beta \)-model,

\[
n_p n_e = n_0^2 \frac{(r/r_c)^{-\alpha}}{(1 + r^2/r_c^2)^{3\beta-\alpha/2}} \frac{1}{(1 + r^2/r_c^2)^{\gamma/2}},
\]

(Vikhlinin et al. 2006). The model accounts for the cuspy centers of clusters as well as the steepening of the profile seen at larger radii. We fix \( \gamma = 3 \) in the above expression (see Vikhlinin et al. 2006).

For the XMM-Newton data, the projected model is convolved with a model of the XMM point-spread function (PSF), as described in Ghizzardi et al. (2001), prior to fitting. The model is normalized by integrating equation 7 over a cylindrical volume with a radius of \( r_{500} \), truncated at 3\( r_{500} \) along the line-of-sight. This is compared to the spectral normalization of a MEKAL model with parameters derived in §3.3 within \( r_{500} \). This way the model normalization, \( n_0 \), can be calculated using the angular distance \( D_A \) determined from the redshift listed in Table 1. The gas mass \( M_{g,500} \) is calculated by setting \( n_e = Z n_p \) and \( \rho_g = m_p n_e A/Z \) where \( Z = 1.199 \) and \( A = 1.397 \) are the average nuclear charge and mass respectively for a plasma with 0.3 solar abundances (implying a mean molecular weight per free electron \( n_e = 1.165 \) assuming abundances by Anders & Grevesse (1989). The gas density \( \rho_g \) is integrated over a sphere of radius \( r_{500} \) determined iteratively for \( M_g \) and \( Y_X \) estimates using \( M_{500} - M_g \) and \( M_{500} - Y_X \) relations, respectively.

The parameter space of the model is explored via Markov chain Monte-Carlo iteration and the uncertainties on the parameters are estimated using the Markov chain posterior.

#### 3.4.1. Cluster images

X-ray images in the 0.5 - 2.0 keV band are generated using the filtered event files and binned into 4" × 4" pixels. These are shown in Figures 5-19 alongside optical images of the clusters from High et al. (2010). Images are smoothed with a 8" Gaussian filter. The images are shown with SZ S/N contours from V10 and the white crosses show the location of the brightest cluster galaxy (BCG). In all cases, the SPT detection is within 1′ of the X-ray centroid and in a majority of cases (9 of 15) it is within 20′. Additionally, Sydney University Molonglo Sky Survey (SUMSS, Mauch et al. 2003) radio sources are marked in the images with small circles (15″ radius). Extended X-ray structures are marked with larger circles (30″ radius) or arrows and are discussed for each individual cluster in the Appendix.

#### 3.5. ICM temperature, gas mass, luminosity and \( Y_X \)

The average spectral temperature of the intra-cluster medium is measured within the [0.15-1]\( r_{500} \) aperture as described in §3.3. A new value of \( r_{500} \) is then estimated using this temperature with equation 1 and a new spectrum is extracted using the new radius. The process is repeated until the value of \( T_X \) has converged. For 5 of

\[\text{http://cxc.harvard.edu/contrib/maxim/acidbg}\]
the clusters, the number of source photons within \([0.15-1]\) \(r_{500}\) is less than 1000 and the signal-to-noise is less than 20. This leads to large uncertainties in the spectral fits and could potentially cause systematic biases in the temperature. For these objects, we instead use the estimate of \(T_X\) within \([0.15-0.5]\) \(r_{500}\), where the signal-to-noise is higher and extrapolate using the fitted relation between the two temperatures from Vikhlinin et al. (2009a):

\[
T_{X,(0.15-1)r_{500}}/T_{X,(0.15-0.5)r_{500}} = 0.9075 + 0.0625 T_{X,(0.15-0.5)r_{500}}.
\]

We also add a 10% uncertainty to these temperature estimates to account for the uncertainty in this relation. The clusters for which this procedure was performed are marked in Table 2. For SPT-CL J0551-5709, a cluster which is spatially coincident with the local cluster Abell S0552, we also perform this procedure, but with a different motivation. The surface brightness to the south and north-west of this cluster suggests that there is a significant amount of flux from cooler gas associated with the Abell cluster. This is also seen in the temperature estimates. \(T_X\) drops from 4.4 ± 0.7 keV to 3.3 ± 0.4 keV when comparing the inner and outer apertures. We include SPT-CL J0551-5709 in the scaling relation fits using the corrected temperature and mark is with a red square in the scaling relation plots.

Similarly, the gas mass within \(r_{500}\), \(M_g,500\), is estimated as described in §3.4 using an initial estimate of \(r_{500}\) from equation 1 and the spectrally derived temperature, \(T_X\). The estimate of \(r_{500}\) is then revised using the implied gas mass fraction and equation 4, and the process is repeated until the gas mass converges.

In this work, we adopt redshift evolution of \(f_g,500\) in the same way as it was applied in the work of Vikhlinin et al. (2009a). The redshift dependence of the gas-mass fraction is not well constrained observationally to high-\(z\) and independent measurements of the mass are needed to study this further. The \(f_g,500\)-\(M\) relation is used here only to determine the radius \(r_{500}\) within which to estimate \(M_g\) and has no impact on the \(Y_X\)-based mass estimates. We do not quote any total masses based on the \(f_g,500\)-\(M\) relation in this work.

SPT-CL J0516-5430 has an unusually extended morphology compared to other objects in the sample and the assumption of spherical symmetry is very approximate. The line-of-sight dimension for this cluster and the inferred gas mass are likely to be overestimated for this reason (see, e.g. Nagai et al. 2007). To keep the analysis analogous to Vikhlinin et al. (2009a) we do not attempt to correct for this here. SPT-CL J0516-5430 is marked with a blue triangle in the scaling relation plots.

When estimating \(Y_X\), we use equation 5 to determine \(r_{500}\) and re-calculate both \(T_X\) and \(M_g\) in a similar way. Cluster simulations show that \(Y_X\) exhibits less scatter with mass compared to \(M_g\) and \(T_X\) individually due to the typical anti-correlation of deviations from the mean of these two observables (Kravtsov et al. 2006). This makes the \(Y_X\) mass estimator less dependent on the dynamical state of clusters, as far as it can be estimated from simulations. The X-ray observables are listed in Table 2.

For some of the clusters in the sample, the X-ray data are not deep enough to detect the X-ray emission with high significance out to \(r_{500}\). For five of the clusters, the signal-to-noise in our radial bins (equation 6) drops below 3 at a radius of about 0.6 \(r_{500}\). These are the same five objects with low signal-to-noise spectra described above. The gas density profile is then constrained primarily by the central surface brightness profile and extrapolated out to \(r_{500}\). This leads to systematic uncertainties that are difficult to account for as they depend on cluster morphology. In our sample we find a typical density slope \(\rho_g \propto r^{-\alpha}\) of \(\alpha = -1.90 \pm 0.30\) at 0.6 \(r_{500}\), where we quote the mean and standard deviation. Varying the slope outside of this radius by ±0.30 typically changes the mass outside 0.6 \(r_{500}\) by 8%. This is also an over-estimate of the uncertainty since density profiles generally become steeper with increasing radius out to \(r_{500}\). To be conservative, we add an uncertainty of 10% on the total gas mass for clusters with poor signal-to-noise outside 0.6 \(r_{500}\).

4. TOTAL MASS ESTIMATES

Using the mass proxies \(T_X\) and \(Y_X\), we estimate the gravitational mass, \(M_{500}\) for the clusters using equations 1 and 5. The mass estimates are listed in Table 3. Clusters are identified as mergers based on their morphology. Clusters with secondary maxima, filamentary structure or significant isophotal centroid shifts are classified as unrelaxed (e.g., Molnar et al. 1993). As discussed in Kravtsov et al. (2006), the \(T_X\) based masses should be multiplied by a factor 1.17 for clusters identified as mergers. Following Vikhlinin et al. (2009a) we correct our masses upward with this factor and add an uncertainty of 5% on the masses (Table 3). Our merger classification is listed in Table 1.

In Table 3, we also give the SZ-derived mass estimates from V10. In V10, the SPT cluster survey was used to constrain cosmological parameters while simultaneously fitting a cluster detection significance-mass relation. The significance-mass relation had priors imposed on its normalization, slope, and redshift evolution that were motivated by simulated thermal SZ maps of the sky. These maps were generated from large dark matter simulations (Shaw et al. 2009) that used a semi-analytic gas model of Bode et al. (2007) which was calibrated such that the simulated clusters matched observed X-ray scaling relations for low redshift (\(z < 0.25\) ) clusters. The maximum likelihood significance-mass relation was then used to generate mass estimates for each cluster, including both statistical and systematic uncertainties. These systematic uncertainties were equivalent to a ∼15% uncertainty on the mass estimate for each cluster.

Comparing the X-ray and SZ-based mass estimates, we find that the SZ-derived masses are lower by a factor of 0.83 ± 0.06, where we use only the statistical uncertainties to quote the uncertainty on the average ratio from a fit to a scaling relation between the two with the slope fixed to 1. Two clusters are 2σ inconsistent with unity and both have larger X-ray mass estimates; SPT-CL J0516-5430 and SPT-CL J0546-5345. V10 notes that clusters at \(z \lesssim 0.3\) will have mass estimates biased low because the power-law scaling that is assumed for the significance-mass relation does not fully capture the effects of CMB-confusion on the SZ signal. This effect could possibly explain the relatively low SZ-inferred mass found for SPT-CL J0516-5430, which is the only cluster in this sample at \(z < 0.3\). Leaving this cluster out, the
average ratio of the SZ- to the X-ray-mass is 0.89 ± 0.06. Overall, we consider the agreement between the X-ray and SZ mass estimates reasonable given the ∼15% systematic uncertainty on the mass estimates quoted in V10. However, there is some evidence that the SZ masses in V10 are biased low by ∼10%.

This difference could have several explanations. For example, the semi-analytic gas model used to calibrate the simulations in V10 could have a redshift evolution that differs from self-similar evolution. There could also be differences in the gas profiles at large radii that cause systematic differences between the measured $Y_X$-based mass and the derived mass from the $M$-$\xi$ scaling in V10 that were not included in the simulations. It should also be noted that the $M_{500}$-$Y_X$ relation used here has been calibrated using X-ray hydrostatic mass estimates.

Numerical simulations have suggested that hydrostatic mass estimates could be biased low by ∼10% due to additional non-thermal pressure support of the gas in clusters (e.g., Lau et al. 2009; Meneghetti et al. 2010). However, comparisons of hydrostatic and weak lensing mass estimates of low redshift ($z < 0.3$) clusters show agreement at the <9% level and suggest no significant bias (e.g., Zhang et al. 2008; Vikhlinin et al. 2009a; Zhang et al. 2010). Additionally, the $M_{500}$-$Y_X$ relation assumes self-similar redshift scaling. The details of redshift-dependent departures from this scaling are not constrained observationally, although the results from simulations provide a 5% upper limit for any evolution of the amplitude of this relation between $z = 0$ and 0.6.

Regardless of the exact reasons, if the SZ mass estimates are biased low, there would be consequences for the cosmological constraints in V10. It would generally favor a lower normalization of the SZ significance-mass relation and hence a larger value of $\sigma_8$; however, these results should still be within the 68% confidence region in V10 due to the priors imposed on the parameters defining the SZ significance-mass relation which effectively include the 15% systematic uncertainty for the masses. The cosmological implications of including the X-ray mass estimates with the V10 results will be the subject of a future paper.

5. SZ DATA ANALYSIS

5.1. Calculating $Y_{SZ}$

The SPT observations, data reduction, and mapping used in this work are identical to V10, and are outlined in §2.1. The analysis differs from V10 at the point where the SPT maps are spatially filtered to identify the cluster candidates. In our case, the cluster candidates have already been identified, and we instead want to calculate the SZ-inferred integrated Comptonization, $Y_{SZ}$, of each cluster. We calculate $Y_{SZ}$ from the SPT 150 GHz maps by spatially filtering them using a filter motivated by the X-ray measurements of each cluster.

In V10, cluster candidates are identified by spatially filtering the SPT maps with a matched filter (Haehnelt & Tegmark 1996; Melin et al. 2006). This is done in the Fourier domain by multiplying the map by

$$\psi(k_x, k_y) = \frac{B(k_x, k_y) S(|\vec{k}|)}{B(k_x, k_y)^2 N_{astro}(|\vec{k}|) + N_{noise}(k_x, k_y)}$$

where $\psi$ is the matched filter, $B$ is the response of the SPT instrument after time-stream processing to signals on the sky, $S$ is the assumed source template, and the noise power has been broken into astrophysical ($N_{astro}$) and noise ($N_{noise}$) components. The noise covariance matrix $N_{noise}$ includes contributions from atmospheric and instrumental noise, while $N_{astro}$ includes power from primary and lensed CMB fluctuations, an unresolved SZ background, and unresolved point sources. In V10, the source template, $S$, was constructed from a projected spherical $\beta$-model.
In this work, we instead use a SZ source template motivated from X-ray measurements of each cluster. The SZ brightness is proportional to the line-of-sight integral of electron pressure, or density times temperature. The profile is assumed to match the product of the best-fit gas density profile to the X-ray measurements of each cluster (equation 7), and the universal temperature profile of Vikhlinin et al. (2006) (equation 12). These profiles are multiplied together to give the radial pressure profile, and projected onto the sky by doing a line-of-sight integral through the cluster. The radial pressure profile is truncated at 3 × r500, where r500 is estimated using equation 5 and given in Table 3. In constructing the spatial filter, we only need to know the spatial shape of the source, so the SZ model is normalized to unity.

For each of the fifteen clusters, we construct a spatial filter using equation 9 and a source model, as described above. The spatially filtered SPT maps are a measure of the normalization of each source model at the cluster location. Using the SPT maps alone to determine the cluster location would bias the YSZ measurements high because of the freedom to maximize the SPT significance by position. Therefore we use the X-ray measured position as a prior on the cluster location. We define the uncertainty in the SPT normalization of the source model as the standard deviation of the spatially filtered map within a 90' band in declination around the cluster. The SPT maps are calibrated in units of K_{CMB}, the equivalent CMB temperature fluctuation required to produce the same power fluctuation. The SPT normalization is converted to Comptonization using the relation

$$\Delta T = y T_{CMB} \left( \frac{e^{x + 1}}{e^x - 1} - 4 \right) [1 + \delta(x, T_e)],$$

where ΔT is the measured decrement in units of K_{CMB}, y is the Comptonization, T_{CMB} is the CMB blackbody temperature of 2.725 K, x = hν/kT_{CMB}, and δ(x, T_e) accounts for relativistic corrections to the SZ spectrum (Itoh et al. 1998; Nozawa et al. 2000). For the frequency dependent terms in equation 10, we calculate their SPT band averaged value, which would effectively correspond to an observing frequency, ν, of 152.9 GHz. The average SPT band is measured from Fourier Transform Spectroscopy measurements of a sample including more than 90% of the SPT 150 GHz detectors.

For the relativistic correction factor, δ(x, T_e), we assume an electron temperature of T_X,r500/1.11, where T_X,r500 is given in Table 2 and the factor of 1.11 is the average ratio of the X-ray spectroscopic and mass-weighted temperature measured in Vikhlinin et al. (2006). The relativistic correction factor is only a weak function of temperature, between 3 to 10 keV it varies by ~ 4%, and we expect this assumption to negligibly affect our results.

The integrated Comptonization, Y_{SZ,cyl}, is calculated for each cluster by integrating its source model over solid angle, Y_{SZ,cyl} = ∫ y(θ)dΩ, normalized to the best-fit SPT Comptonization, y. To more easily compare to the X-ray measurements, we convert our measurements to units of M_⊙/eV, and define Y_{SZ,cyl} as,

$$Y_{SZ,cyl} = \left( \frac{\mu_e m_p}{\sigma_T} \right) D_A^2 \int y(\theta)d\Omega,$$

where D_A is the angular distance to the source, σ_T is the Thomson cross-section, m_e is the electron mass, m_p is the proton mass, c is the speed of light, μ_e is the mean molecular weight per free electron. In Table 4, we give Y_{SZ,cyl} integrated out to the angular radius corresponding to r500 as determined from the X-ray measurements and given in Table 3. The uncertainty in Y_{SZ} is calculated as the quadrature sum of the uncertainty in the SPT calibration, the measured SPT normalization of the source model, and an additional uncertainty in the assumed source model, which will be discussed in §5.3.

5.2. Spherical Deprojection

Both the X-ray and SZ observations are measuring a projected signal that is proportional to the integrated

<table>
<thead>
<tr>
<th>Name</th>
<th>z</th>
<th>r500^a</th>
<th>M500,Y_X</th>
<th>M500,T_X</th>
<th>M500,SZ,Y_{10}^b</th>
</tr>
</thead>
<tbody>
<tr>
<td>J0000-5748</td>
<td>0.74</td>
<td>950</td>
<td>5.32 ± 1.16</td>
<td>6.74 ± 5.14</td>
<td>2.89 ± 0.61 ± 0.41</td>
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<tr>
<td>J0009-5342</td>
<td>0.4626</td>
<td>1026</td>
<td>5.43 ± 0.60</td>
<td>6.71 ± 3.40</td>
<td>4.26 ± 0.74 ± 0.60</td>
</tr>
<tr>
<td>J0156-5430</td>
<td>0.2952</td>
<td>1463</td>
<td>11.84 ± 1.25</td>
<td>12.34 ± 2.13</td>
<td>6.48 ± 0.95 ± 1.13</td>
</tr>
<tr>
<td>J0528-5300</td>
<td>0.7648</td>
<td>775</td>
<td>2.97 ± 0.89</td>
<td>3.05 ± 4.00</td>
<td>2.83 ± 0.60 ± 0.38</td>
</tr>
<tr>
<td>J0533-5005</td>
<td>0.8810</td>
<td>656</td>
<td>2.06 ± 0.53</td>
<td>2.25 ± 1.71</td>
<td>2.71 ± 0.56 ± 0.37</td>
</tr>
<tr>
<td>J0546-5345</td>
<td>1.0665</td>
<td>840</td>
<td>5.33 ± 0.62</td>
<td>5.25 ± 1.13</td>
<td>3.25 ± 0.51 ± 0.44</td>
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<tr>
<td>J2331-5051</td>
<td>0.5707</td>
<td>972</td>
<td>4.70 ± 0.51</td>
<td>4.91 ± 1.36</td>
<td>4.63 ± 0.73 ± 0.66</td>
</tr>
<tr>
<td>J2332-5358</td>
<td>0.32</td>
<td>1134</td>
<td>5.66 ± 0.48</td>
<td>6.69 ± 1.08</td>
<td>5.19 ± 0.85 ± 0.83</td>
</tr>
<tr>
<td>J2337-5942</td>
<td>0.7814</td>
<td>1046</td>
<td>7.43 ± 0.90</td>
<td>8.10 ± 2.18</td>
<td>6.32 ± 0.84 ± 0.97</td>
</tr>
<tr>
<td>J2341-5119</td>
<td>0.9983</td>
<td>847</td>
<td>5.06 ± 0.66</td>
<td>5.20 ± 1.69</td>
<td>4.05 ± 0.58 ± 0.63</td>
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<tr>
<td>J2342-5411</td>
<td>1.08</td>
<td>647</td>
<td>2.47 ± 0.32</td>
<td>2.39 ± 0.63</td>
<td>2.65 ± 0.50 ± 0.37</td>
</tr>
<tr>
<td>J2355-5056</td>
<td>0.35</td>
<td>1014</td>
<td>4.18 ± 0.43</td>
<td>4.26 ± 1.57</td>
<td>4.17 ± 0.80 ± 0.63</td>
</tr>
<tr>
<td>J2359-5009</td>
<td>0.76</td>
<td>816</td>
<td>3.45 ± 0.67</td>
<td>4.97 ± 2.61</td>
<td>3.32 ± 0.60 ± 0.46</td>
</tr>
</tbody>
</table>

Note. — Masses estimated from X-ray mass-proxy relations in equations 1 and 5. SZ derived masses from V10 are shown for comparison.

^a Estimated using the M_{500} – Y X relation (equation 5).

^b Quoted masses from V10 include statistical and systematic uncertainties and have been scaled with our adopted h.
gas properties in a line of sight through the cluster. This project has different physical dependencies between X-ray and SZ observations that must be considered. To ease this comparison we deproject each measurement so that they correspond to a physical observable that is proportional to a spherical volume integral of each signal.

The $Y_X$ estimates are deprojected as described in §3.4. For X-rays, the effect of the deprojection is smaller than for the SZ, because the X-ray flux is proportional to $n_e^2$ and only weakly dependent on temperature. This effect decreases the contribution to the X-ray signal from large radii where the density is lower. Also, $Y_X$ will be proportional to a X-ray spectroscopic weighted temperature from gas between 0.15$r_{500}$ and $r_{500}$, while $Y_{SZ}$ is related to the gas mass weighted temperature. We do not formally account for this difference in calculating either $Y_X$ or $Y_{SZ}$. However later in this section we will discuss this effect in interpreting their comparison.

We calculate $Y_{SZ,sph}$ by spherically deprojecting the $Y_{SZ,cyl}$ measurements in Table 4. The SZ signal is proportional to the line of sight integral of the electron density, $n_e$, times temperature, $T_e$. For each cluster, we assume the density profile derived from the X-ray imaging analysis and assume a temperature profile of a universal form,

$$T(r) = T_0 \frac{(x/0.045)^{1.9} + 0.45}{(x/0.045)^{1.9} + 1} \frac{1}{(1 + (x/0.6)^2)^{0.45}},$$

(12)

where $x = r/r_{500}$ (see Vikhlinin et al. 2006). We then define $Y_{SZ,sph}$ as:

$$Y_{SZ,sph} = \frac{Y_{SZ,cyl}}{C} = \frac{1}{C} \int_{r_{500}}^{r_{500}} \frac{T(r)n_e(r)dr}{\int_{r_{500}}^{r_{500}} T(r)n_e(r)dr},$$

(13)

where $C$ is the ratio of the integrals of pressure in a cylindrical and spherical volume through the cluster. For these integrals we use an integration radius of $r_{500}$ and truncate the density and temperature radial models at 3$r_{500}$. We note that varying the truncation radius from between 3-5$r_{500}$ changes our measurement of $Y_{SZ}$ in Table 4 by less than 1%. For the $r_{500}$ aperture, the median and standard deviation of $C$ across the sample is 1.23 ± 0.08. This is consistent with the value of 1.203 that is calculated assuming the universal pressure profile from Arnaud et al. (2009).

As noted earlier in this section, even after the above spherical deprojection, $Y_X$ and $Y_{SZ,sph}$ are not directly comparable because of their different weighting of the electron temperature. $Y_X$ is proportional to the X-ray spectroscopic weighted temperature, $T_X$, while $Y_{SZ,sph}$ is proportional to the mass weighted temperature, $T_{mg}$. The size of this difference has been estimated by several authors from X-ray measurements. In Vikhlinin et al. (2006), they estimate $T_X/T_{mg} = 1.11$ for a sample of relaxed massive clusters with high angular resolution X-ray temperature measurements between 70 kpc and $r_{500}$. This would imply that $Y_{SZ,sph}/Y_X = 1.203$. In Arnaud et al. (2009), a similar analysis was performed for a sample that included both relaxed and un-relaxed clusters, and it was estimated that $Y_{SZ,sph}/Y_X = 0.924 ± 0.004$. For the measured density profiles and assumed temperature profiles for our sample, we would expect an average ratio of $Y_{SZ,sph}/Y_X = 0.91 ± 0.01$. To compute this expected ratio we have used the “spectroscopic-like” temperature (e.g., Mazzotta et al. 2004), given the density and temperature profiles used here.

In §6.2 we compare the above ratios, predicted purely from X-ray observations, to the ratio including the $Y_{SZ,sph}$ calculated from the SZ data as described in §5.1 and deprojected as described in this section.

### 5.3. Model Uncertainty

The integrated Comptonization, $Y_{SZ}$, inferred from the SZ data depends on the assumed spatial model through the construction of the spatial filter, and the volume integral through the deprojection the SZ data. For a typical cluster, the X-ray data constrains the cluster density profile with high signal-to-noise out to $\sim r_{500}/2$ with no information on the temperature profile. Since a significant amount of the SZ signal is coming from larger radii than this, we would like to estimate how much uncertainty our assumed profile is adding to the $Y_{SZ}$ estimates. To help do this, we calculate the $Y_{SZ}$ of each cluster assuming the universal pressure profile of Arnaud et al. (2009), variations of the temperature profile in equation 12, and an isothermal model.

Arnaud et al. (2009) measure a universal pressure profile from X-ray measurements for a representative set of local massive clusters. These measurements were deep enough to constrain the cluster density and temperature profiles out to $r_{500}$ in each cluster. They find that their sample is well fit by a universal pressure profile that is defined only by $M_{500}$. For each cluster in our sample, we re-calculate $Y_{SZ,sph}$ assuming the universal pressure profile from Arnaud et al. (2009) and the X-ray measured values for $r_{500}$ and $M_{500}$ given in Tables 4 and 3, respectively. Comparing these $Y_{SZ,sph}$ estimates to the values given in Table 4, we find that this ratio has a mean and standard deviation of 1.01 ± 0.15 averaged over the sample. There is one significant (> 20%) outlier, SPT-CL J0516-5430, whose $Y_{SZ,sph}$ deviates by a factor of 0.59. This cluster is a major merger, and the observed cluster profile is noticeably poorly fit by the Arnaud et al. (2009) pressure profile, which does not capture the disturbed distribution of the central gas in this cluster. Leaving this cluster out of our comparison, the average ratio becomes 1.03 ± 0.10. Therefore we see no detectable bias when assuming the Arnaud et al. (2009) pressure profile, but these results suggest that there could be an addi-
tional $\sim 10\%$ uncertainty in our $Y_{SZ,sph}$ measurements from our assumed pressure profile.

As an additional test, we vary the outer slope of the Vikhlinin et al. (2006) temperature profile in equation 12, letting the exponent on the term in the right hand side of the denominator vary between 0.16 to 0.80, from its starting value of 0.45. This range of values matches the full range of effective slopes of the temperature profile at $r_{500}$ as measured in the sample of Vikhlinin et al. (2006). Calculating new $Y_{SZ,sph}$ estimates assuming these new temperature profiles, we find they change the $Y_{SZ,sph}$ values given in Table 4 on average by a factor of 1.09 $\pm$ 0.06 and 0.91 $\pm$ 0.05 for the exponent values of 0.16 and 0.80, respectively. While we consider such a significant systematic shift in the temperature profile unlikely across the whole sample, this level of variation in temperature profiles could contribute added statistical uncertainty in the $Y_{SZ,sph}$ estimates.

As a limiting case, we also recalculate $Y_{SZ,sph}$ assuming that the cluster is isothermal and with a density profile corresponding to the best-fit profile to the X-ray measurements. Comparing $Y_{SZ,sph}$ for the isothermal profile to the values given in Table 4, we find that this ratio has a mean and standard deviation of 1.12 $\pm$ 0.08 averaged over the sample. This is certainly an extreme case because of the abundance of evidence for the temperature profile in clusters decreasing significantly by $r_{500}$ (e.g., Vikhlinin et al. 2006; Arnaud et al. 2009). Recent Suzaku observations of nearby clusters at large radii also indicate a drop of gas temperature by a factor $\sim 3$ at $r_{200}$ (e.g., Fujita et al. 2008; George et al. 2009; Bautz et al. 2009; Reiprich et al. 2009; Hoshino et al. 2010). We also note that the temperature profile with an exponent of 0.80 assumed in the previous paragraph would have a gas temperature that dropped by a factor of $\sim 5$ at $r_{200}$, so even this would seem like an extreme case.

Overall, we are encouraged that the variation in $Y_{SZ,sph}$ is found to be $\sim 10\%$ when assuming a broad range of different temperature and pressure profiles found in the works of Vikhlinin et al. (2006) and Arnaud et al. (2009). Therefore we conservatively assume an additional $10\%$ statistical uncertainty in our $Y_{SZ,sph}$ estimates from any assumed model uncertainty, which we have added in quadrature to the uncertainties given in Table 4.

5.4. Point Source Contamination

Astrophysical point sources in the direction of the cluster can potentially fill in the SZ cluster decrement and bias the integrated Comptonization low. There are generally two populations of sources that can contaminate the SZ signal: dusty or radio bright sources. In V10, the potential contamination from both were discussed, and neither is expected to significantly bias the SZ measurements at 150 GHz averaged over the sample. We review some of those conclusions here.

In the Appendix, we discuss radio detections at 843 MHz by the Sydney University Molonglo Sky Survey (SUMSS) towards the clusters in this work. The majority of radio sources in clusters have been found to have steep spectra with $\alpha < -0.5$ (where $S \propto \nu^\alpha$) (e.g., Coble et al. 2007; Lin et al. 2009). For example, Coble et al. (2007) find a median spectral index of $-0.72$ between 1.4 and 28.5 GHz for radio sources towards a sample of massive clusters ranging from $0.14 < z < 1.0$. In V10, they noted that a typical cluster would suffer a decrease of $\Delta \xi = 1$ for a 2 mJy (5 mJy) source at 150 GHz, located at 0.5' (1') from the cluster center. Assuming a spectral index of $-0.72$, a 2 mJy (5 mJy) source at 150 GHz would be $\sim 83 \,(210) \text{mJy}$ at the SUMSS observing frequency of 843 MHz. As detailed in the Appendix, no source has been detected above either threshold within 1' of any cluster in this work. However, for any individual cluster radio source, extrapolating its flux from radio frequencies to the SPT observing frequency of 152.9 GHz is difficult because of the broad range and frequency dependence of the spectral indices that is typical for these sources (e.g., Coble et al. 2007; Lin et al. 2009). In practice, extrapolating the flux of any SUMSS detected source to the SPT observing frequency will have factors of a few uncertainty without further observations spanning an intermediate range of frequencies. With this caveat in mind, there appears to be no source strong enough to significantly bias the SZ flux estimate for any cluster in this work.

Including every source detected within 0.5' of each cluster center in our sample, the cumulative SUMSS flux is 224.5 mJy. If we assume a spectral index of -0.72, this would imply an average flux of $\sim 0.35 \text{mJy per cluster}$ at 152.9 GHz. This would correspond to an average decrease in the SZ flux of $\sim 2\%$, assuming that these sources represent an overdensity to the background population.

The average radio source contamination in SZ surveys has been recently estimated in simulations of the microwave sky by Sehgal et al. (2010). These simulations were motivated by observations by Lin et al. (2009) measuring the radio source population characteristics in low redshift ($z < 0.25$) X-ray detected clusters. For the mass limit of this work, $M_{500} \approx 3 \times 10^{14}h^{-1}M_\odot$, these studies predict that in $\lesssim 1\%$ of clusters there would be radio source contamination large enough to effect the SZ flux measurement at the $> 20\%$ level, with this result largely independent of redshift. V10 estimates a similar rate of contamination using a radio source count model (de Zotti et al. 2005) and the measured over-abundance of radio sources near clusters (Coble et al. 2007). Overall, the combination of the above results lead us to not expect any significant radio source contamination.

Emission from sub-millimeter bright galaxies can also contaminate the SZ signal. These sources are typically dusty star forming galaxies that are very luminous but highly obscured, such that their luminosity is peaked at infrared (IR) wavelengths. The number counts and fluxes of these sources have been characterized between 0.5 - 2.0 mm wavelength by several experiments (e.g., Coppin et al. 2006; Devlin et al. 2009; Vieira et al. 2009; Austermann et al. 2010), and these measurements reasonably match simple analytic models describing their source population distribution (Lima et al. 2010b). These sources can largely be approximated as a random background that contributes additional signal in the SPT maps, and we have explicitly accounted for them in our matched filter in equation 9. However, this implementation did not account for any emission that could be correlated with clusters, or additional noise from the background objects being gravitationally lensed by the cluster. The former was argued to be insignificant in V10 because the IR overdensity towards clusters is expected to be $\ll 1 \text{mJy}$ at 150 GHz even out to highest redshift...
clusters in our sample, and observational measurements of the 100 \( \mu \)m flux towards local clusters (\( z \sim 0.2 \)) are too low to significantly bias the SZ flux measurements even allowing for a strong redshift evolution of these sources. The latter has been characterized by Lima et al. (2010a) to increase the flux noise towards clusters by \( \sim 60\% \), which is at a level such that it would be \( \lesssim 3\% \) of the SZ flux for the mass range of clusters in our sample. We note that because lensing is a surface brightness conserving process, this latter effect would not bias our results when averaged over a large sample of clusters.

One noteworthy exception where point source contamination is evident is the cluster SPT-CL J2332-5358. This cluster was the only candidate in V10 that was also coincident with a point source detected in the SPT 220 GHz maps at \( > 5\sigma \). Similar 220 GHz detected point sources have a source flux, \( S \), that scales with frequency as \( S \propto \nu^\alpha \) with \( \alpha = 3.3 \pm 0.7 \) between 150 and 220 GHz (Vieira et al. 2009). Extrapolating its measured 220 GHz source flux would predict a flux of \( 11.0 \pm 2.7 \) mJy at 150 GHz. To the 150 GHz SPT maps, we subtract a point source of this brightness, after convolving it with the 150 GHz SPT transfer function. We then repeat the same cluster extraction method as in V10 to calculate a new cluster position and SPT-significance, \( \xi \), which is given in Table 1. This point source corrected 150 GHz map is used when calculating the \( Y_{\text{SZ}} \) given in Table 4, and the added uncertainty in \( Y_{\text{SZ}} \) from the uncertain point source spectral index is added in quadrature with the other uncertainties outlined in §5.1.

6. SCALING RELATIONS

Here we discuss the X-ray and SZ scaling relations for the sample. We consider the integrated SZ flux \( Y_{\text{SZ}} \) for both cylindrical and spherical volumes of integration. The integrated cylindrical estimate is deprojected to a spherical estimate as described in §5.2.

To measure the slope and normalization of the scaling relations we perform linear regression in the presence of intrinsic scatter. We use the Bayesian method of Kelly (2007) and maximize the probability of a linear model in log-log space, accounting for measurement uncertainties in both axes. For the \( Y_{\text{SZ}} \) scaling relations below we fit the relations with another method, accounting for the \( \xi \) selection. We investigate this fitting method using simulated cluster samples as detailed in §6.2. We have assumed self-similar \( E(z) \) scaling for these scaling relations unless otherwise stated. The best fit parameters for these relations are listed in Table 5.

6.1. \( M_g - T_X, L_X - T_X \) and \( L_X-M_{500} \)

The scaling between the ICM gas mass (\( M_g \)) and the spectroscopic temperature (\( T_X \)) is an important test of the properties of this SZ-selected sample. We want to know if the SZ selection biases the observables in any way or if they are consistent with those from X-ray selected samples. The values of \( T_X \) and \( M_g \) within \( r_{500} \) used here were derived as described in §3.3 and §3.4 and are given in Table 2. This relation can provide clues about the high-\( z \) behavior of the \( M-T_X \) and \( M-M_g \) scaling relations (equations 1 and 4). To accurately constrain the evolution of these relations, however, either hydrostatic masses or masses obtained via an independent measurement are needed for high-\( z \) clusters.

We fit the relation

\[
E(z)M_g = 10^A \left( \frac{T_X}{5 \text{ keV}} \right)^B ,
\]

to the data as described above, accounting for intrinsic scatter. Figure 1 shows the \( M_g - T_X \) relation with the best fit power law (solid line).

To compare our results with previous studies using X-ray selected samples we use the gas mass and temperature data for the low-\( z \) (median \( z \approx 0.05 \)) and high-\( z \) (median \( z \approx 0.48 \)) as measured by Vikhlinin et al. (2009a). We choose these samples since the data analysis is analogous to ours. We fit the scaling relation to these samples using the same method. For our SPT-selected sample (median \( z = 0.74 \)) we find \( A = 13.65 \pm 0.10 \) and \( B = 1.95 \pm 0.66 \) (see Table 5). When fitting the low-\( z \) sample, we find \( A = 13.54 \pm 0.02 \) and \( B = 1.66 \pm 0.08 \), at 5 keV. For the high-\( z \) sample from Vikhlinin et al. (2009a), the best fit parameters are \( A = 13.66 \pm 0.03 \) and \( B = 1.64 \pm 0.21 \) in even better agreement with our findings. This offset between low-\( z \) and high-\( z \) samples indicate deviations from self-similarity. The slope appears steep in our relation, although consistent with both the low- and high-\( z \) X-ray selected samples as well as the self-similar slope, \( B = 1.5 \). A positive trend of the gas mass fraction with mass has been observed previously (e.g., Mohr et al. 1999; Vikhlinin et al. 2006) and is likely what is causing the steeper than self- similar slope of the relation.

Numerical simulations (e.g. Kravtsov et al. 2005) also show an increasing gas mass fraction with redshift, presumably due to a different distribution of mass progenitors at high redshift. The SPT-selected cluster sample is consistent with this scenario. The normalization agrees well with that of the high-\( z \) sample of Vikhlinin et al. (2009a) while still consistent with the low-\( z \) sample at the 1\( \sigma \) level. We cannot rule out that this behavior is partly related to deviations from self-similarity in the \( M-T_X \) relation at high redshift. However, the results are consistent with previous findings from X-ray selected cluster samples.

We also investigate the luminosity-temperature relation (e.g., Edge & Stewart 1991; Markevitch 1998) for the sample where we compare the luminosity in the 0.5-2.0 keV band within \( r_{500} \), including the core, \( L_X \), to the temperature estimate within \( r_{500} \), \( T_X \), with the core excised. We fit the relation

\[
E(z)^{-1}L_X = 10^A \left( \frac{T_X}{5 \text{ keV}} \right)^B ,
\]

and find a shallow slope, \( B = 1.92 \pm 0.60 \) (see Table 5), marginally inconsistent with most previous works which typically find a slope close to 3 (e.g. Mantz et al. 2009). For completeness, we also estimate the bolometric luminosity by extrapolating the spectrum for all frequencies using the core excised temperature. This introduces additional uncertainties on the luminosity due to the uncertainty in the temperature estimate. For the \( L_{\text{bol}}-T_X \) relation, we find a steeper slope of \( B = 2.31 \pm 0.85 \), consistent with previous work. The best fit relation is shown in Figure 2 (left).

We investigate the luminosity-mass relation (e.g., Reiprich & Böhringer 2002) and adopt the best-fit \( E(z) \) scaling found by Vikhlinin et al. (2009a) to fit for the
relation

\[ E(z)^{-1.85} L_X = 10^A \left( \frac{M_{500, Y_X}}{10^{15} M_\odot} \right)^B. \] (16)

The best fit relation is shown in Figure 2 (solid line) compared to the best fit in Vikhlinin et al. (2009a) (dashed line). We find a slope of \( B = 1.16 \pm 0.20 \), shallow compared to the \( B = 1.61 \pm 0.14 \) found by Vikhlinin et al. (2009a), while the normalization agrees very well near the median mass of the sample at \( 5 \times 10^{14} M_\odot \). From the relation in Vikhlinin et al. (2009a) we expect a normalization of \( A = 44.25 \) at \( 5 \times 10^{14} M_\odot \), which is in good agreement with our measured value of \( A = 44.28 \pm 0.07 \). The slope of the fit is mostly driven by SPT-CL J0516-5430 (blue triangle) which has a high derived mass for its measured luminosity, possibly caused by the observed merging activity in the cluster. Excluding this object from the fit changes the slope to \( B = 1.45 \pm 0.29 \), consistent with Vikhlinin et al. (2009a).

In summary, we find a good agreement between the X-ray observables measured for this SZ-selected sample when comparing our results to X-ray selected samples where similar analysis methods were used.

6.2. \( Y_{SZ} \) scaling relations

6.2.1. Fitting method

The selection of clusters above a fixed S/N threshold can bias the SZ-X-ray scaling relations. The SZ flux, \( Y_{SZ} \), is correlated with the SZ S/N on which the sample is selected and ignoring this effect will lead to bias in the \( Y_{SZ} \)-scaling relation parameters. The steepness of the cluster mass function makes this bias more of a problem since the number of clusters in a given mass bin grows exponentially towards lower masses. This causes more clusters to be near the selection threshold where the effect is most prominent. The measurement uncertainty in \( Y_{SZ} \) and the intrinsic scatter in the scaling relation cause clusters with low mass (or \( Y_X \)) to scatter over the selection threshold. This is visible as a tail towards low mass (or \( Y_X \)) near the \( Y_{SZ} \)-selection threshold in the scaling relation plots. If ignored, this bias will lead to a biased slope and a higher normalization near the threshold in the fitted scaling relation parameters. Below we describe our method to account for the SZ selection in our scaling relation fits.

We estimate how the selection cut in \( \xi \) translates to a selection in \( Y_{SZ, sph} \) using simulated SPT observations. We repeat the procedure used in V10 to estimate the SZ selection function and SZ-significance mass relation, where observations of simulated SZ maps are added to the dominant sources of astrophysical confusion and instrumental noise, mock observed, and processed through the SPT analysis pipeline. In contrast from V10, we keep track of the predicted value of \( Y_{SZ, sph} \) for each cluster, which is estimated using the \( Y_{sph} - M_{500} \) scaling relation from Arnaud et al. (2009) and the simulated cluster’s mass, \( M_{500} \). Also, we measure \( \xi \) for each simulated cluster in the same way as we do for the SPT observations, where we record the maximum significance S/N over different spatial filters and in a single realization of the astrophysical confusion and noise.

With these results we find a best-fit relation

\[ \xi = 5.90 \left( \frac{Y_{SZ}}{2 \times 10^{14} M_\odot \text{keV}} \right)^{0.64} E(z)^{0.16} \] (17)

with intrinsic scatter in log \( Y_{SZ} \) of 0.10. Comparing this result with the measured \( \xi \) and \( Y_{SZ} \) for our sample we find that the simulation-based relation is offset slightly high in \( \xi \). When fitting the normalization of the above relation with the observed data we find a best fit value of \( 5.56 \pm 0.31 \), about 6% lower than the simulations suggest. The reason for this offset may be related to systematic uncertainties in the simulation of SPT observations. We adopt this best fit normalization from the observed data, and use the above relation to construct a selection function in log \( Y_{SZ} \). We estimate the selection probability as an error function in log \( Y_{SZ} \) at this threshold with the width set by the intrinsic scatter in this relation. To account for the uncertainty in the simulations we use include a 10% gaussian uncertainty on the normalization and a 20% gaussian uncertainty on the intrinsic scatter, which we marginalize over for our results.

The \( Y_{SZ} \) scaling relation fit is performed as follows. First, we calculate the probability of measuring a certain \( Y_{SZ} \) for a given \( M \) or \( Y_X \) and a given scaling relation with log-normal scatter. This is then convolved with the
measurement uncertainty in $Y_{SZ}$. The $Y_{SZ}$ selection cut is applied by multiplying the probability of measuring $Y_{SZ}$ with an error function in log $Y_{SZ}$ as described above. The likelihood is then calculated as the product of probabilities for all clusters and is maximized to obtain the scaling relation parameters.

The uncertainties on the calculated $Y_{SZ}$ are correlated with the measurement error in $Y_X$ and $M_{500,YX}$. The $r_{500}$ radius used for integration of $Y_{SZ}$ is estimated from $M_{500,YX}$ and any measurement scatter in $Y_X$ will contribute to scatter in $Y_{SZ}$ of similar magnitude. This will lead to scatter along the scaling relation and for this reason we ignore these effects. We fit for the intrinsic scatter in the $Y_{SZ}$-$M_{500,YX}$ and $Y_{SZ}$-$Y_X$ relations while noting that the scatter in $Y_{SZ}$ and $Y_X$ with the true cluster mass will also be partially correlated.

To test our fitting method for potential biases we run it on mock samples drawn from a fiducial mass function. We perform this test both in the limit of a sample with 1000 clusters and for many samples of 15 clusters. To estimate a $Y_{SZ}$ for the mock samples we assume a $Y_{SZ}$-$M$ relation with log-normal scatter and convolve this with a linear measurement scatter of $0.4 \times 10^{14} M_\odot$. We assume that we know the relation between $\xi$ and $Y_{SZ}$ and its scatter to reproduce the sample selection function. In the limit of a large sample of clusters, we find no measurable bias in either the fitted slope or normalization. For 1000 generated mock samples of 15 clusters we find that the scatter in normalization and slope between the samples and the average measurement error on the parameters over the samples are consistent. The normalization is reproduced to within 0.1% of the input value while we find that the slope is biased low by 2.0%. The average measurement error on the slope from the 15 cluster samples is 13.5% so this bias is of little significance.

6.2.2. $Y_{SZ}$-$Y_X$ relation

The relationship between $Y_{SZ}$ and $Y_X$ is determined by details of the gas and temperature distribution in the cluster. It is effectively measuring a relationship between the mass weighted and X-ray spectroscopic weighted temperature. In Arnaud et al. (2009), X-ray measurements of both relaxed and unrelaxed clusters were used to approximate a $Y_{SZ}$-$Y_X$ relation, for a spherical integration to $r_{500}$. Their work found a relationship with a normalization of $0.924 \pm 0.004$, implying a lower mass weighted temperature. This result is consistent with previous X-ray measurements (e.g., Vikhlinin et al. 2006), and our expectation of $0.91 \pm 0.01$ for the measured gas density profiles and the assumed temperature profile in this work. Hydrodynamical simulations predict a similar ratio, but also find that X-ray measurements could overestimate the integrated pressure at some level (e.g., Nagai et al. 2007). This would imply that the X-ray predictions of the normalization of the $Y_{SZ}$-$Y_X$ relation could be biased high; for example, the results of Nagai et al. (2007) suggest a bias at the few percent level.

Recently there has been some evidence that the $SZ$ signal from clusters could be lower than this expectation from comparison of WMAP SZ observations of X-ray selected clusters (Komatsu et al. 2010), however similar analyses have led to different conclusions (Melin et al. 2010). By fitting a normalization of the $Y_{SZ}$-$Y_X$ relation to the clusters in this work, we can test for a similar offset, which could be indicative of differences in the real gas profiles from those we’ve assumed or some other systematic bias in either the $SZ$ or X-ray measurement.

Using the method described in §6.2.1, we fit a scaling relation between $Y_{SZ}$ and $Y_X$ of the form

$$Y_{SZ} = 10^A \left( \frac{Y_X}{3 \times 10^{14} M_\odot \text{keV}} \right)^B.$$  \hspace{1cm} (18)

In Figure 3, we show the $Y_{SZ}$-$Y_X$ relation for both cylindrical (left) and spherical (right) $Y_X$, denoted by $Y_{SZ,cyl}$ and $Y_{SZ,sph}$ respectively. In Table 5, we give fits to equation 18 given several different assumptions, which are described and discussed further below.

For both cylindrical and spherical $Y_{SZ}$ we find a slope consistent with unity. The dashed line in the plots is the $Y_{SZ} = Y_X$ relation and the dotted line (right) shows the expected $Y_{SZ,sph}/Y_X$ ratio of 0.924 found in the work of Arnaud et al. (2009). We also fit the scaling relations with the slope fixed to 1 (see Table 5). The normalization of our $Y_{SZ,sph}/Y_X$ fit implies an average ratio of $0.82 \pm 0.07$, consistent with the prediction of Arnaud et al. (2009) at the 1.5$\sigma$ level. The most significant outlier from
the $Y_{SZ,sph}$-$Y_X$ relation, SPT-CL J0000-5748, also has the lowest $Y_{SZ,sph}/Y_X$ ratio in our sample. If we exclude this cluster from the fit, we find an even better agreement with a normalization of 0.87±0.06. The marginally lower normalization found in this relation when compared to the X-ray prediction is expected at some level from hydrodynamical simulations which predict that $Y_{SZ}$ could be biased high when estimated using X-ray data.

As an additional test, we fit the $Y_{SZ}$-$Y_X$ relation including a factor of $E(z)^F$, and fit for $C$ while keeping the slope fixed to 1. We find $C = -0.36 ± 0.58$, consistent with no evolution. Further investigating this, we divide the sample in two redshift bins. The low-$z$ bin consists of the 7 clusters at $z < 0.7$ while the high-$z$ bin consists of the 8 clusters at $z > 0.7$. Fixing the slope to 1, we find a ratio of 0.88 ± 0.12 for the low-$z$ sample, and a ratio of 0.72 ± 0.14 for the high-$z$ sample. If we exclude the most significant outlier from the high-$z$ sample, SPT-CL J0000-5748, we find a ratio of 0.83 ± 0.09.

As discussed at the beginning of this section, the $Y_{SZ,sph}/Y_X$ ratio is expected to be lower than 1 due to the different weighting of the temperature in the two estimates. As explained in §5.3, our measurement of $Y_{SZ}$ is sensitive to the assumed temperature profile, whereas the $Y_X$ measurement is not. This causes the measured $Y_{SZ,sph}/Y_X$ ratio to change under different assumptions about the temperature profile. However, the expected ratio of the gas-mass weighted and the spectroscopic weighted temperatures changes by a similar amount as the temperature profile changes. The overall effect makes the measured ratio to be fairly independent to our assumptions of the shape of the temperature profile. This is particularly true for variations in the slope of the temperature profile at large radii that we tested in §5.3.

Overall we find a normalization consistent with our expectations given the gas profiles in this work, and believe this result to be largely independent of our assumed temperature profile. We find no significant evidence for a redshift dependent evolution in the normalization of the $Y_{SZ}$-$Y_X$ relation, but note that the high-$z$ sub-sample does favor a $\sim 1$σ lower normalization. This could be marginal evidence for redshift evolution, however further X-ray and SZ observations will be needed to make any statistically significant statements.

6.2.3. $Y_{SZ}$-$M_{500}$-$Y_X$ relation

Finally, we investigate the relation between the SZ flux, $Y_{SZ}$, and the X-ray derived mass, $M_{500}$. This is not an independent result from the previous section because the X-ray masses are calculated directly from the $Y_X$ measurements and the $M_{500}$-$Y_X$ relation in Vikhlinin et al. (2009a). However, it is useful in understanding the SZ mass calibration, and can be compared to previous measurements of this relation.

Using the method of §6.2.1, we fit a scaling relation between $Y_{SZ}$ and $M_{500}Y_X$ of the form

$$Y_{SZ} = 10^A \left( \frac{M_{500}Y_X}{3 \times 10^{14} M_\odot} \right)^B E(z)^{2/3}. \quad (19)$$

In Figure 4 we show the $Y_{SZ} - M_{500}Y_X$ relation, and in Table 5 we give the best-fit parameters to equation 19. The slope of this relation is found to be 1.67 ± 0.29 for $Y_{SZ,sph}$-$M_{500}Y_X$ with intrinsic scatter of 0.09 ± 0.05 (see Table 5). This slope is consistent with the self-similar expectation of 5/3 and previous measurements (e.g., Bonamente et al. 2008; Arnaud et al. 2009). In Arnaud et al. (2009), a slope of 1.790 ± 0.015 is measured for their predicted $Y_{SZ,sph}$ using X-ray observables.

We also fit the relation keeping the slope fixed at the value expected from X-ray studies, 1.79, and note that the normalization of the $Y_{SZ}$-$M$ relation is lower than that of Arnaud et al. (2009). The Arnaud et al. (2009) results imply a normalization $A = 14.115 ± 0.003$, using our adopted $h$ while we find $A = 14.03 ± 0.04$, around a 2σ offset. The best fit relation is shown in Figure 4 (solid line) with the Arnaud et al. (2009) relation shown (dashed line) for comparison.

It should be noted that part of the offset is due to differences in the mass estimates in our work and in the work of Arnaud et al. (2009). The normalization in the $M - Y_X$ relation (equation 5), used here to estimate cluster masses, is different from the normalization of the $M - Y_X$ relation used in Arnaud et al. (2009) to derive the $Y_{SZ}$-$M_{500}$ relation in that work. For our adopted $h$, we find a mass, $M_{500,Y_X} = 4.83 ± 0.17 \times 10^{14} M_\odot$, at $Y_X = 3 \times 10^{14} M_\odot$ keV from equation 5. If we instead use the $M - Y_X$ relation from Arnaud et al. (2009), we find $M_{500,Y_X} = 4.64 ± 0.12 \times 10^{14} M_\odot$, at the same $Y_X$. To account for this scaling relation offset we shift the normalization of the Arnaud et al. (2009) $Y_{SZ}$-$M_{500}$-$Y_X$ relation down by a factor $(4.83/4.64)^{1.79} = 1.074$, resulting in $A = 14.084 ± 0.003$. This is reduces the offset to our measurement, $A = 14.03 ± 0.04$, to a similar level to what we found for the $Y_{SZ}$-$Y_X$ relation in the previous section.

It is important to note that the intrinsic scatter in this scaling relation does not directly reflect the low scatter relationship between $Y_{SZ}$ and the gravitational mass. The mass is derived directly from $Y_X$ and therefore includes the scatter in the $Y_X$-mass relationship. Independent measurements of the gravitational mass, e.g., through weak lensing or hydrostatic masses from X-ray data, are necessary to diagnose scatter in the $Y_{SZ}$-$M$ relation.

7. CONCLUSIONS

We have presented the results from X-ray observations of a subset of 15 clusters from the first SZ selected cluster catalog from the South Pole Telescope cluster survey. We report the X-ray properties of this sample, including measurements of $M_\mu$, $T_X$, and $Y_X$, and have used $T_X$ and $Y_X$ to estimate the total masses of the clusters. We find generally good agreement between the expected X-ray properties of this sample and those expected from scaling relations. However, we find some indication of deviations from self-similar evolution of the $M_\mu$-$T$ relation compared to other local cluster samples at a level consistent with the explanation of an evolving gas mass fraction in high redshift clusters.

Using the X-ray measured cluster positions and gas profiles, we have re-analyzed the SZ measurements to calculate each cluster’s integrated Compton-y parameter, $Y_{SZ}$. We further use the X-ray measured gas profiles to deproject the SZ measurements so that they correspond to a spherical integrated Comptonization, $Y_{SZ,sph}$, that is more directly comparable to the X-ray measurements.

We have calculated scaling relations between $Y_{SZ,sph}$ and the $Y_{SZ,sph}$-$Y_X$ relation, SPT-CL J0000-5748, also has the lowest $Y_{SZ,sph}/Y_X$ ratio in our sample. If we exclude this cluster from the fit, we find an even better agreement with a normalization of 0.87±0.06. The marginally lower normalization found in this relation when compared to the X-ray prediction is expected at some level from hydrodynamical simulations which predict that $Y_{SZ}$ could be biased high when estimated using X-ray data.

As an additional test, we fit the $Y_{SZ}$-$Y_X$ relation including a factor of $E(z)^F$, and fit for $C$ while keeping the slope fixed to 1. We find $C = -0.36 ± 0.58$, consistent with no evolution. Further investigating this, we divide the sample in two redshift bins. The low-$z$ bin consists of the 7 clusters at $z < 0.7$ while the high-$z$ bin consists of the 8 clusters at $z > 0.7$. Fixing the slope to 1, we find a ratio of 0.88 ± 0.12 for the low-$z$ sample, and a ratio of 0.72 ± 0.14 for the high-$z$ sample. If we exclude the most significant outlier from the high-$z$ sample, SPT-CL J0000-5748, we find a ratio of 0.83 ± 0.09.

As discussed at the beginning of this section, the $Y_{SZ,sph}/Y_X$ ratio is expected to be lower than 1 due to the different weighting of the temperature in the two estimates. As explained in §5.3, our measurement of $Y_{SZ}$ is sensitive to the assumed temperature profile, whereas the $Y_X$ measurement is not. This causes the measured $Y_{SZ,sph}/Y_X$ ratio to change under different assumptions about the temperature profile. However, the expected ratio of the gas-mass weighted and the spectroscopic weighted temperatures changes by a similar amount as the temperature profile changes. The overall effect makes the measured ratio to be fairly independent to our assumptions of the shape of the temperature profile. This is particularly true for variations in the slope of the temperature profile at large radii that we tested in §5.3.

Overall we find a normalization consistent with our expectations given the gas profiles in this work, and believe this result to be largely independent of our assumed temperature profile. We find no significant evidence for a redshift dependent evolution in the normalization of the $Y_{SZ}$-$Y_X$ relation, but note that the high-$z$ sub-sample does favor a $\sim 1$σ lower normalization. This could be marginal evidence for redshift evolution, however further X-ray and SZ observations will be needed to make any statistically significant statements.

6.2.3. $Y_{SZ}$-$M_{500}$-$Y_X$ relation

Finally, we investigate the relation between the SZ flux, $Y_{SZ}$, and the X-ray derived mass, $M_{500}$. This is not an independent result from the previous section because the X-ray masses are calculated directly from the $Y_X$ measurements and the $M_{500}$-$Y_X$ relation in Vikhlinin et al. (2009a). However, it is useful in understanding the SZ mass calibration, and can be compared to previous measurements of this relation.

Using the method of §6.2.1, we fit a scaling relation between $Y_{SZ}$ and $M_{500}Y_X$ of the form

$$Y_{SZ} = 10^A \left( \frac{M_{500}Y_X}{3 \times 10^{14} M_\odot} \right)^B E(z)^{2/3}. \quad (19)$$

In Figure 4 we show the $Y_{SZ} - M_{500}Y_X$ relation, and in Table 5 we give the best-fit parameters to equation 19. The slope of this relation is found to be 1.67 ± 0.29 for $Y_{SZ,sph}$-$M_{500}Y_X$ with intrinsic scatter of 0.09 ± 0.05 (see
and the X-ray measured quantities $Y_X$ and $M_{500,YX}$. We fit the $Y_{SZ,sph},Y_X$ relation, and find a slope consistent with unity, $0.96 \pm 0.18$. Fixing this slope to 1, we re-fit the relation and find a normalization that implies a ratio of $Y_{SZ,sph}/Y_X = 0.82 \pm 0.07$. This normalization effectively corresponds to the ratio between the mass weighted and X-ray spectroscopic weighted temperature. For the spherically symmetric density and temperature profiles assumed in this work, we would have expected a ratio of $0.91 \pm 0.01$, consistent with predictions from other X-ray studies of clusters (see e.g. Vikhlinin et al. 2006; Arnaud et al. 2009). We therefore find a normalization of the $Y_{SZ,sph},Y_X$ relation that is marginally consistent with, although lower than this prediction. This indicates that the SZ and X-ray measured pressure largely agree. We find, however, that the lower normalization is more pronounced in the $z > 0.7$ high-$z$ sub-sample, with a ratio of $0.72 \pm 0.14$. Further X-ray and SZ observations of high redshift clusters are needed to determine if this is a real effect.

Using the $Y_X$ measurement as a proxy for the total cluster mass with a relation calibrated in Vikhlinin et al. (2009a), we find similar results when we fit a $Y_{SZ,sph} - M_{500,YX}$ relation. We find a slope consistent with the self-similar expectation of $Y_{SZ} \propto M^{5/3}$ and a normalization consistent with the predictions from X-ray measurements by Arnaud et al. (2009). We have compared the $Y_X$ inferred total mass to the SZ significance-inferred total mass from Vanderlinde et al. (2010). Considering only the clusters used in their cosmological analysis, we find an average ratio of $M_{500,SZ}$ to $M_{500,YX}$ of $0.89 \pm 0.06$, which is within the $\sim 15\%$ systematic uncertainty on the simulation-based mass estimates in Vanderlinde et al. (2010).

This work is encouraging for future studies of SZ-selected clusters. It suggests that SZ mass estimates are consistent with X-ray based mass estimates used in other cluster cosmology studies. This result is important for the use of SZ-selected cluster samples to constrain cosmology, and demonstrates that the X-ray measurements can play a valuable role in calibrating SZ surveys. This work also highlights the potential power of SZ surveys to study cluster evolution due to the broad redshift range of SZ-selected samples. The SPT has now surveyed an additional $\sim 900$ deg$^2$ to a similar depth to the $\sim 178$ deg$^2$ used in this work, and expects to cover over 2000 deg$^2$
by December 2011. These measurements will produce a catalog of hundreds of massive SZ-discovered clusters that extend to high redshift. Joint X-ray and SZ measurements of these clusters promise to place interesting constraints on cluster formation and gas physics for the most massive young clusters in the universe.

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Dressler, A., Hare, T., Bigelow, B. C., & Osip, D. J. 2006, Proc. SPIE, 6269,

TABLE 5

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<thead>
<tr>
<th>Relation</th>
<th>$L_X/T_X$</th>
<th>$L_{bol}/T_X$</th>
<th>$L_X/M_{500}$</th>
<th>$M_p/T_X$</th>
<th>$Y_{SZ,sph}-Y_X$</th>
<th>$Y_{SZ,cyl}-Y_X$</th>
<th>$Y_{SZ,sph}-Y_X$</th>
<th>$Y_{SZ,cyl}-Y_X$</th>
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<tr>
<td></td>
<td>44.16 ± 0.09</td>
<td>1.92 ± 0.60</td>
<td>1.16 ± 0.20</td>
<td>1.85</td>
<td>0.11 ± 0.06</td>
<td>5 keV</td>
<td>0.17 ± 0.08</td>
<td>5 keV</td>
<td>0.11 ± 0.04</td>
<td>$10^{15} M_\odot$</td>
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<td>44.64 ± 0.14</td>
<td>2.31 ± 0.85</td>
<td>1.16 ± 0.20</td>
<td>1.85</td>
<td>0.11 ± 0.04</td>
<td>$10^{15} M_\odot$</td>
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<tr>
<td></td>
<td>13.65 ± 0.10</td>
<td>1.95 ± 0.66</td>
<td>−1</td>
<td>0.11 ± 0.06</td>
<td>5 keV</td>
<td>0.09 ± 0.04</td>
<td>$3 \times 10^{14} M_\odot$ keV</td>
<td></td>
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<tr>
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<td>14.40 ± 0.08</td>
<td>0.96 ± 0.18</td>
<td>0</td>
<td>0.09 ± 0.04</td>
<td>$3 \times 10^{14} M_\odot$ keV</td>
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<td>14.50 ± 0.05</td>
<td>0.90 ± 0.17</td>
<td>0</td>
<td>0.09 ± 0.04</td>
<td>$3 \times 10^{14} M_\odot$ keV</td>
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<td></td>
<td>14.60 ± 0.10</td>
<td>1.67 ± 0.29</td>
<td>2/3</td>
<td>0.09 ± 0.04</td>
<td>$3 \times 10^{14} M_\odot$ keV</td>
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<td>14.03 ± 0.04</td>
<td>1.79</td>
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Note. — Self-similar $E(z)$-scaling has been assumed here except for the $L_X/M_{500}$ relation, where the best fit evolution from (Vikhlinin et al. 2009a) is adopted. $\sigma_y$ represents the intrinsic scatter in log $y$.
In this section we discuss the X-ray emission for each of the objects in the sample and note any extended substructures or other features in the X-ray images. Nearby radio sources, as listed in the SUMSS survey are quoted with fluxes at 843 MHz, and we also discuss any particularly bright X-ray point-sources. The X-ray images ($\sim 15' \times 15'$) in the 0.5-2.0 keV band are shown in Figures 5-19 with optical $grz$-images from the Magellan Baade 6.5-m telescope (see High et al. 2010). For SPT-CL J0546-5345 and SPT-CL J2342-5411 we instead show $grz$-images from the BCS survey since Magellan images were not available in all filters for these objects. Both X-ray and optical images are overlaid with SZ significance, $\xi$, contours from V10. The contours are from spatially filtered SPT maps where the clusters appear as positive significance detections, and where the spatial filtering often causes noticeable negative significance ringing around the cluster. The contour levels are spaced at $1.5\sigma$ increments starting from zero with dashed contours for $\xi \leq 0$ and solid contours for $\xi > 0$. The position of the BCG (Stalder et al., in prep.) is marked with a white cross in the X-ray images. Large circles and arrows show the locations of interesting X-ray features, small circles correspond to locations of SUMSS sources mentioned in the text. North is up and east is left.

**SPT-CL J0000-5748**

This cluster is the highest redshift ($z = 0.74$) cluster in the sample with a sharp central brightness peak (see Figure 5). The X-ray peak is coincident with the BCG and is also associated with a SUMSS radio source with a flux of $40.4 \pm 1.5$ mJy indicative of a central AGN. The central brightness peak is broader than the Chandra PSF and we do not associate the central X-ray emission with AGN emission.

**SPT-CL J0509-5342**

There is a clear double-peak in the central X-ray emission of this cluster (see Figure 6). The main peak is associated with the central BCG. The secondary peak is located $\sim 20''$ west of the main component (arrow) and the X-ray brightness ratio is 1/4. The secondary peak is also associated with two elliptical galaxies. The peak of the SZ flux is located between the two components.

The main component is the most centrally compact core in this sample which indicates that a possible merger is in an early stage. A bright SUMSS radio source with a flux of 125.7 mJy is located 2.8' WSW of the cluster and is also coincident with an X-ray point source. This is not likely associated with the cluster. A possible third cluster component shows up as an extended X-ray source 3.5' NNW of the cluster (large circle) and is coincident with two elliptical galaxies in the optical image.

**SPT-CL J0516-5430**

This apparently merging cluster is very elongated along the N-S direction and also features a string of galaxies with the same alignment (see arrow, Figure 7). The BCG is also offset by 30'' from both the X-ray and SZ peaks. There is a second extended X-ray component 9.5' SW of the main cluster (SW arrow). This possible subcluster also shows up in the SZ map as a S/N> 2 detection. The subcluster is outside of the field of the Magellan observation and no galaxy appears to be associated with this component in a Digitized Sky Survey image.

**SPT-CL J0528-5300**

This is a faint cluster ($L_X(0.5 - 2.0 kev) = 1.9 \pm 0.2 \times 10^{44}$ erg s$^{-1}$) without much structure and the BCG is offset by 15'' from the X-ray peak (Figure 8). The BCG is coincident with a SUMSS radio source with a flux of 61.2 $\pm$ 2.0 mJy likely associated with a central AGN.

**SPT-CL J0533-5005**

This is another faint cluster ($L_X(0.5 - 2.0 kev) = 1.2 \pm 0.3 \times 10^{44}$ erg s$^{-1}$) with no well defined core. The BCG is offset 45'' from the central X-ray emission indicating an unrelaxed state (Figure 9). The SZ and X-ray peaks are also offset by 30''. A local galaxy at $z = 0.0147$ is located 3' NW of the cluster (circle) but is not likely to affect X-ray or SZ measurements.

**SPT-CL J0546-5345**

The X-ray image in Figure 10 shows a substructure extending $\sim 1'$ SW from the main cluster (arrow) suggesting that a minor merger may be taking place. This elongation is aligned with an apparent extension of the SZ signal further supporting this scenario. One of the X-ray point sources, 6' NW, is associated with a SUMSS radio source with a radio flux of 19.1 $\pm$ 0.9 mJy.

**SPT-CL J0551-5709**

This apparently merging cluster has a disturbed X-ray morphology and the X-ray and SZ peaks are offset by 30''. A SUMSS radio source with flux of 22.7 $\pm$ 1.6 mJy is located 15'' east of the BCG in a region with an overdensity of cluster galaxies (Figure 11). Another radio source is located 5' NE of the cluster and is associated with an X-ray source and a local galaxy.
The X-ray morphology and temperature structure of this object suggest that it is projected on top of a low redshift structure. Extended emission can be seen both towards the S and NW, covering a large area beyond $r_{500}$. The temperature decreases from 4.4 keV to 3.3 keV when increasing the aperture radius from $0.5 r_{500}$ to $r_{500}$, implying that this extended emission has very low temperature. The emission is likely associated with Abell S0552, located at this position. High et al. (2010) find a strong red sequence at $z = 0.09$ at the cluster location. The X-ray emission associated with this local cluster will likely bias any measurement of the temperature due to the contribution of the low-$z$ flux at large radii. We attempt to correct for this bias as described in §3.5.

**SPT-CL J0559-5249**

The peak of the X-ray emission is offset to the south by $\sim 25''$ with respect to the BCG location and SZ peak in this cluster (Figure 12). The X-ray emission south of the BCG is associated with a group of galaxies at this location. The BCG is coincident with a $50.1 \pm 2.7$ mJy SUMSS radio source. On larger scales, the X-ray emission and SZ signal are extended to the SW, clearly visible $2.5''$ from the cluster core (arrow in Figure 12). This cluster is clearly in a merging state.

**SPT-CL J2331-5051**

As noted in V10, this cluster appears to be a pair of merging clusters with the smaller component located $2.5'\text{ SE}$ of the main cluster. A spectral analysis of the smaller component indicates that it has a luminosity of $2 \times 10^{43}$ erg s$^{-1}$, roughly 5% of the main cluster. As can be seen in Figure 13, the location of the X-ray emission and SZ decrement at the subcluster show very good agreement. This subcluster was also confirmed (see High et al. 2010) to be at the same redshift as the main cluster. Additionally, this cluster has a well defined X-ray cavity $18''$ S of the X-ray peak (arrow). This cluster also hosts a cool core, coincident with the BCG and a $13.3 \pm 1.0$ mJy SUMSS radio source.

**SPT-CL J2332-5358**

This nearby ($z = 0.32$) regular cluster shows no sign of significant structure in our XMM-Newton observation with 7 ks of good data. The location of the BCG, the X-ray peak and the peak of the SZ-signal show good agreement and the cluster is likely relaxed (see Figure 14). This cluster is also the only candidate in V10 with a point source detected in the SPT 220 GHz maps at $> 5\sigma$. We correct the SZ flux and position of the cluster by subtracting this source from the 150 GHz assuming a dust spectral index as described in §5.4.

**SPT-CL J2337-5942**

This cluster shows a sharp X-ray surface brightness discontinuity NE of the cluster center (arrow in Figure 15). Interestingly, the BCG is also located in this region, providing tentative evidence that this is a compact gas core moving through the ambient cluster medium, driving a shock front. The cluster shows an overall irregular morphology without a clear peak in the X-ray emission. This is the cluster with the highest SZ signal in the sample and is the second most massive with $M_{500, \text{X}} = 7.43 \pm 0.80 \times 10^{14} M_\odot$. There are two SUMSS radio sources with $7.1 \pm 1.0$ mJy and $7.8 \pm 1.0$ mJy fluxes $33''$ and $23''$ respectively from the cluster center.

**SPT-CL J2341-5119**

In this cluster there are no signs of significant structure within the cluster in the X-ray data (see Figure 16). However, $2.5'\text{ NW}$ of the cluster center (large circle) there is a faint extended X-ray source at the location of a bright galaxy indicating that this is a satellite galaxy group about to merge with the main cluster. The BCG shows good correspondence with the X-ray peak and the main cluster appears relaxed. There is a $7.9 \pm 1.1$ mJy SUMSS radio source $6''$ N of the BCG which could be associated with a central AGN.

**SPT-CL J2342-5411**

This high redshift ($z = 1.08$) cluster shows no signs of significant merger activity and is likely relaxed. However there is a possible secondary component or tail to the SW, $15''$ from the core.

**SPT-CL J2343-5521**

In a 70 ks Chandra observation, no X-ray source was found at the location of this SPT detection. This field was also imaged with BCS and Magellan with no optical counterpart found. The optical data suggested that this cluster would have to be at $z \gtrsim 1.2$ to not find an optical counterpart in either observation (High et al. 2010). The long Chandra observation allows us to put strict upper limits on the luminosity. Assuming the cluster is at $z = 1$, the $5\sigma$ upper limit is $L_X(0.5-2.0\text{keV}) < 3 \times 10^{42}\text{ergs s}^{-1}$. The same limit assuming $z = 2$ is $10^{43}$ ergs s$^{-1}$. Additionally, no cluster counterparts could be identified in a 3.6 $\mu$m Spitzer/IRAC observation. The X-ray, IR, and optical data all strongly support that this is a single band false detection by SPT.

**SPT-CL J2355-5056**

This cluster shows good agreement between the position of the BCG and the X-ray peak and shows no significant structure in the X-ray image (Figure 18). The SZ peak is offset $\sim 35''$ to the south of the BCG while there is no evidence for extended X-ray emission in this direction. The X-ray emission is peaked in the center, likely due to a cool core.
The X-ray image in Figure 19 reveals a 45'' offset between the location of the SPT SZ detection and the X-ray peak. The BCG position is offset from the SZ peak by 33''. Located east of the cluster is a local (z = 0.029) galaxy pair (circle, 7.5' E) which is clearly seen in both the X-ray and optical images. Another local galaxy (z = 0.047) is associated with a faint X-ray source (1.8' W), this is excluded from our X-ray analysis. A very bright point source is also visible in the X-ray image (3.5' NW), likely associated with an AGN. We do not find a counterpart to this source in the NASA/IPAC extragalactic database (NED) and estimate a flux of $1.6 \times 10^{-13}$ erg s$^{-1}$ cm$^{-2}$ in the 0.5-2.0 keV band. A SUMSS radio source is located 15'' SE of the BCG with a flux of $21 \pm 1.1$ mJy. The X-ray morphology is elongated along the east-west direction, possibly indicating merging activity. This is also supported by the ~70 kpc offset between the X-ray peak and the BCG position.
Fig. 7.— SPT-CL J0516-5430, z=0.2952

Fig. 8.— SPT-CL J0528-5300, z=0.7648
Fig. 9.— SPT-CL J0533-5005, z=0.8810

Fig. 10.— SPT-CL J0546-5345, z=1.0665
Fig. 11.— SPT-CL J0551-5709, z=0.4230

Fig. 12.— SPT-CL J0559-5249, z=0.6112
Fig. 13.— SPT-CL J2331-5051, $z=0.5707$

Fig. 14.— SPT-CL J2332-5358, $z=0.32$
Fig. 15.— SPT-CL J2337-5942, z=0.7814

Fig. 16.— SPT-CL J2341-5119, z=0.9983
Fig. 17.— SPT-CL J2342-5411, z=1.08

Fig. 18.— SPT-CL J2355-5056, z=0.35
Fig. 19.— SPT-CL J2359-5009, z=0.76
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