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# Consolidation effects on uplift capacity of shallow horizontal plate anchors in dilating sand

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# Consolidation effects on uplift capacity of shallow horizontal plate anchors in dilating sand

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## Abstract

This paper examines the effect of consolidation on a shallowly embedded horizontal plate anchor in medium dense and dense sand using centrifuge tests, where the loading rate was varied over four orders of magnitude. The experimental results show a 4.8 to 5.5 times increase in anchor capacity as the consolidation condition changes from drained to undrained, driven by a steadily increasing negative excess pore pressure with increasing loading rate. At the highest loading rate, under undrained conditions, the measured maximum negative excess pore pressures reaches a steady limit suggesting the occurrence of cavitation. This increase in anchor capacity with non-dimensional displacement rate is captured using a ‘backbone curve’ interpretation framework that describes the change in capacity between the limiting values of drained and undrained anchor capacity. Calculation of drained anchor capacity is straightforward relative to the more challenging problem of calculating undrained capacity, particularly during cavitation. This was addressed separately through a numerical parametric study (pure undrained analysis) using a bounding surface soil model involving different water depths (cavitation potential), densities and embedment ratios. The numerical results are then systematised into a simple extended analytical solution to allow estimation of undrained anchor capacity under different densities and water depths.

## Keywords

Plate anchors; consolidation; cavitation; bounding surface model; sand

## List of notations

$\delta$	Anchor displacement
$\delta_{soil}$	Soil displacement
$\sigma_{vo}$	Total overburden stress
$\sigma'_v$	Vertical effective stress

1	$\sigma'_n$	Normal effective stress
2	$\mu$	Kinematic viscosity
3	$\mu_{20}$	Kinematic viscosity at 20° C
4		
5	$\Delta\gamma_{oct}^p$	Incremental octahedral plastic shear strain
6		
7	$\gamma'$	Effective unit weight of soil
8	$\phi$	Friction angle
9		
10	$\psi_d$	Dilation angle
11	$\tau_f$	Shear stress
12		
13	$\theta$	Lode angle
14		
15	$\psi$	State parameter
16	$A$	Area of the plate anchor
17	$B$	Plate width (strip or rectangular anchor)
18		
19	$c_v$	Consolidation coefficient
20		
21	$d$	Dimension
22	$d_{50}$	Median particle size
23		
24	$e$	Void ratio
25	$F_{u,ps,dr}$	Uplift factor in plane strain under drained conditions
26		
27	$g$	Centrifuge acceleration level
28	$H$	Embedment depth of the anchor plate
29		
30	$I_R$	Relative dilatancy index
31		
32	$\alpha, \beta, j$	Fitting constants
33	$K_0$	Coefficient of earth pressure
34	$K_w$	Bulk modulus of water
35		
36	$L$	Plate length (rectangular anchor)
37		
38	$N_\gamma$	Anchor factor
39	$N_{\gamma,dr}$	Drained anchor factor
40		
41	$N_{\gamma,un}$	Undrained anchor factor
42	$P_a$	Atmospheric pressure
43		
44	$p'$	Mean effective stress
45		
46	$q_{net}$	Net cone resistance
47	$q$	Mobilised anchor capacity
48		
49	$q_u$	Peak anchor capacity, resistance
50	$q_{u,dr}$	Peak anchor capacity (drained conditions)
51		
52	$q_{u,un}$	Peak anchor capacity (undrained conditions)
53		
54	$q_{u,ref}$	Resistance at a reference velocity
55	$R_D$	Relative density
56	$S_f$	Shape factor
57		
58	$T$	Temperature
59	$t$	Thickness of plate anchor
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1	$V$	Dimensionless velocity
2	$V_{dr}$	Drained limit for dimensionless velocity
3	$V_{un}$	Undrained limit for dimensionless velocity
4		
5	$V_{50}$	Dimensionless velocity at 50% degree of consolidation
6		
7	$u_{cav}$	Cavitation pressure of water (negative value)
8	$u_{hydro}$	Initial hydrostatic pore pressure at soil surface
9		
10	$\Delta u_{max}$	Maximum magnitude of pore pressure change limited by cavitation
11	$\Delta u_{cav}$	maximum change of pore pressure possible at a location prior to cavitation
12		
13	$u$	Pore pressure
14	$u_2$	Measured pore pressure at cone shoulder
15		
16	$v$	Velocity
17	$z$	Depth
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# 1 Introduction

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4 2 Increasing motivation to extract energy from offshore renewable sources has seen the development of a  
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6 3 number of floating renewable energy devices for harnessing wind and wave power (Hallowell et al. 2018, Al  
7  
8 4 Shami et al. 2019). Such devices are subjected to high uplift and/or lateral loads under wind and wave action, such  
9  
10 5 that a safe and reliable anchoring system for keeping these devices on station is a critical component of the overall  
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12 6 infrastructure. An efficient mooring solution for these applications could be provided by plate anchors due to their  
13  
14 7 high efficiency in resisting uplift loading (Bradshaw et al. 2015, O’Loughlin et al. 2018, Al Hakeem & Aubeny  
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16 8 2019). To explore the mooring potential of plate anchors for offshore renewable energy applications, recent studies  
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18 9 on plate anchors have focused on investigating their performance in sands, which are typical of seabed conditions  
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20 10 where floating renewable energy devices might be located (e.g. Knappett et al. 2015, Giampa et al. 2017, Chow  
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22 11 et al. 2018).

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25 12 Installation (and anchoring) of floating energy devices are relevant for relatively deep sea, where they can  
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27 13 take advantage of superior wind speeds and high wave energy (e.g. the Hywind floating wind farm moored in 95  
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29 14 to 129 m water depths off the coast of Scotland). Under such aggressive conditions, the anchors will be subjected  
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31 15 to rapid loading, leading to changing consolidation conditions or ‘rate effects’ in soil resistance and hence in  
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33 16 anchor capacity. Rate effects on soil resistance comprise two components: (a) consolidation effects due to  
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35 17 generation of excess pore pressure, and (b) viscous (rheological) effects (Lehane et al. 2009). For the problem  
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37 18 considered here (in sand), the contribution of viscous effects is negligible relative to consolidation effects (e.g.  
38  
39 19 Dayal & Allen 1975, Watanabe & Kusakabe 2013), such that in the current context ‘rate effects’ are considered  
40  
41 20 analogous to ‘consolidation effects’. The significance of consolidation effects during rapid shearing in sand was  
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43 21 recognised as early as the 1970s in studies of retrieval lift operations of objects embedded in the seafloor (Muga  
44  
45 22 1968); and later in various other offshore applications, such as ploughing (van Os & van Leussen 1987, Palmer  
46  
47 23 1999), in-service performance of suction caissons for offshore wind turbines (Achmus & Thieken 2014, Bienen  
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49 24 et al. 2018), upheaval buckling of pipelines (Bransby & Ireland 2009, Williams et al. 2013) and installation of  
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51 25 drag embedded fluke anchors (Heurlin et al. 2015). In particular, consolidation effects could have severe  
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53 26 consequences in dense, dilating sand (as typically found offshore), where dilation induced negative excess pore  
54  
55 27 pressures result in significantly increased soil resistance during ploughing or anchor installation (e.g. Palmer 1999,  
56  
57 28 Heurlin et al. 2015).

Quantification of consolidation effects for boundary valued problems can be achieved using a dimensionless velocity term ( $V = vd/c_v$ ), where  $v$ ,  $d$  and  $c_v$  denote velocity or loading rate, drainage path length and the coefficient of consolidation for the soil, respectively (Finnie & Randolph 1994). With increasing  $V$ , the normalised soil resistance for a given boundary value problem (i.e. resistance at a given velocity over the reference undrained/drained resistance at a reference velocity  $q_u/q_{u,ref}$ ) varies non-linearly as the consolidation regimes changes from drained to partially drained to undrained, with  $V_{dr}$  and  $V_{un}$  marking the boundary between drained and undrained regimes respectively (e.g. Finnie & Randolph 1994, Chung et al. 2006, Lee & Randolph 2011). For applications in sand (Fig. 1), this variation in resistance, commonly known as a ‘backbone curve’, can be quantified using (Chow et al. 2020):

$$\frac{q_u}{q_{u,dr}} = \frac{1 + \frac{q_{u,un}}{q_{u,dr}} \left(\frac{V}{V_{50}}\right)^c}{1 + \left(\frac{V}{V_{50}}\right)^c} \quad (1)$$

where  $q_{u,un}/q_{u,dr}$  is the ratio of the limiting undrained resistance to the reference drained resistance,  $V_{50}$  is the value of  $V$  at 50% degree of consolidation and  $c$  is a constant controlling the evolution of  $q_u/q_{u,dr}$  within the partially drained region. Note that the original expression in Chow et al. (2020) also considered viscous rate effects, which are ignored here as the change in resistance due to viscous effects will be negligible relative to the change due to consolidation effects (as discussed earlier).

Chow et al. (2020) also highlights the challenges in estimating undrained capacity of an anchor in sand analytically, particularly in the scenario where cavitation may occur in a dilating dense sand. Cavitation occurs when vapour bubbles form within the pore water at an absolute magnitude of pore water pressure equal to the vapour pressure of water (i.e. approximately -100 kPa). The potential for cavitation to occur is influenced by the initial hydrostatic pore pressure ( $u_{hydro}$ ), as evident from triaxial tests (Mcmanus & Davis 1997) and suction bucket tension loading tests (Andersen et al. 1994, Bye et al. 1995, Byrne 2000, Kelly et al. 2003, Bienen et al. 2018). Since cavitation is only triggered at around -100 kPa, the potential for cavitation to occur for a soil in a state of shear (and at a given density) is higher in the water depths relevant for floating renewable energy devices.

This paper investigates consolidation effects for vertical uplift capacity ( $q_u$ ) of a horizontal plate anchor shallowly embedded in sand through centrifuge modelling and numerical analysis. The centrifuge study involved 26 uplift tests on pre-embedded rectangular plate anchors in both medium dense and dense sand, where the uplift velocity varied over four orders of magnitude. In order to extend the consolidation regime from drained to undrained, the coefficient of consolidation ( $c_v$ ) was scaled down in the tests by reducing the soil permeability,

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achieved by saturating the sand using a viscous pore fluid instead of water, similar to the approach described in Adamidis & Madabhushi (2015). Consolidation effects on plate anchor capacity are then quantified using a backbone curve that spans between the limiting drained and undrained capacities. The limiting undrained anchor capacity ( $q_{u,un}$ ) under cavitation is further quantified through numerical analyses under a range of soil conditions and subsequently synthesised in a simple analytical framework, such that it may be used readily as a calculable input for the backbone curve.

## Experimental methodology

### *Modelling using viscous pore fluid*

In order to reduce  $c_v$ , a viscous pore fluid known as methocel cellulose ether (Food grade F450) at a concentration of 2.2% was used to saturate the sand samples. The methocel solution was prepared according to the hot/cold mixing technique (Dow 2002). The viscosity of the solution is shearing rate and temperature dependent (Adamidis & Madabhushi 2015). These dependencies were assessed using a series of rotational vane viscometer tests across a temperature ( $T$ ) range between 17°C and 32°C (Fig. 2) and a range of rotation rates up to 20 s<sup>-1</sup>, which covers the expected range of shearing rates ( $v/d = 0.0003$  to 3.13 s<sup>-1</sup>) for the soil in the anchor uplift tests. With no significant change expected in the measured kinematic viscosities ( $\mu$ ) over the relatively low range of shearing rates considered here (Adamidis & Madabhushi 2015), the variation in with temperature was fitted using the exponential expression:

$$\frac{\mu}{\mu_{20^\circ\text{C}}} = 2.23e^{-0.04T} \quad (2)$$

with  $\mu_{20^\circ\text{C}} = 530.3$  cSt. Temperature in the centrifuge was continuously monitored during testing, so Eqn. (2) allowed accurate estimation of the pore fluid viscosity for evaluating  $c_v$  and  $V$  for the anchor tests.

### *Model plate anchors*

Sand blasted stainless steel rectangular plate anchors of width ( $B$ ) 20 mm, length ( $L$ ) 40 mm and thickness ( $t$ ) 6 mm were used in this study (see Fig. 3(b)). The anchor width was selected to ensure that potential grain size effects are negligible, achieved by ensuring that the ratio of the plate width to the mean particle size ( $B/d_{50}$ ) is greater than 35 (Garnier et al. 2007). Selected plates were instrumented with a pore pressure transducer (PPT) (MS54XX with an absolute pressure range of 700 kPa, with a 1.5 mm thick 35 micron polyethylene filter over the sensing element) on the top face of the plate. The PPTs were placed along the central axis at a distance  $0.5B$  from the centre of the plate. Each PPT was calibrated in a pressure chamber prior to the anchor tests, with essentially

85 identical calibration factors obtained for both positive and negative pressure ranges. As shown in Fig. 3(b), the  
86 plates were attached to a smooth modular rigid rod of diameter 4.7 mm and length 80 mm. The rod extended to  
87 above the sand surface, such that it was accessible to an electrical actuator, creating a means of loading the plate  
88 anchor (see Fig. 3).

#### 89 *Sample preparation details and centrifuge testing procedure*

90 The anchor uplift tests were conducted in a fine sub-angular silica sand, with properties as summarised in  
91 Table 1. The sand samples were prepared in centrifuge sample containers ('strongboxes') at two target densities  
92 ( $R_D$ ) of 70% and 48% by air pluviation using an automated sand rainer. Strongboxes of two different sizes were  
93 used, the first measuring 1300 mm (length)  $\times$  390 mm (breadth)  $\times$  225 mm (height) for the dense sample ( $R_D =$   
94 70%) and the second measuring 650 mm (length)  $\times$  390 mm (breadth)  $\times$  325 mm (height) for the two medium  
95 dense samples ( $R_D \sim 48\%$ ). A 20 mm thick gravel layer overlain by a geotextile was placed at the bottom of the  
96 strongboxes to facilitate drainage (Fig. 3(a)). To avoid consideration of potential (and variable) installation effects  
97 in the interpretation, the anchors were pre-located in the sample (analogous to the 'wished-in-place' approach in  
98 numerical modelling). This was achieved by pausing pluviation when the sand reached the required anchor  
99 embedment depth, vacuum levelling the sand surface and careful placement of the anchor. Prior to placing the  
100 plate anchors, the plate anchor PPTs were saturated in methocel solution for at least 12 hours. Pluviation was then  
101 resumed and the top of the sample vacuum levelled such the anchors were at an embedment depth of  
102 approximately 60 mm, giving an approximate plate embedment ratio ( $H/B$ ) to the top of the plate of 2.7.

103 The sample was then saturated from the base with the methocel solution using a gravity feed method (see  
104 Chow et al. 2020). A free layer of the methocel fluid was maintained over the sample during testing including a 4  
105 mm surface layer of baby oil to prevent evaporation and subsequent viscosity changes of the pore fluid in the  
106 samples. The height of this layer of free fluid was  $\sim 21$  mm and  $\sim 35$  mm for the dense and medium dense samples  
107 respectively, giving a hydrostatic pressure ( $u_{hydro}$ ) of 10.3 kPa and 17.2 kPa respectively at 50g at the soil surface.  
108 To ensure saturation was maintained for the plate anchor PPTs, pluviation and sample saturation was completed  
109 as soon as possible (typically within 10 hours) after placing the anchors. In-flight measurements of pore pressure  
110 (on the anchor) were hydrostatic before loading the anchor, confirming that saturation of the PPTs was maintained.

111 The density of each sample (and hence  $R_D$ ) was calculated from global measurements of the sample mass  
112 and volume prior to locating the anchors. The appropriateness of this 'average' relative density over the sample  
113 height was examined using the empirical correlations developed for the same sand (Roy et. al 2019) from profiles

114 of net cone resistance ( $q_{net}$ ) obtained using a miniature piezocone, 10 mm in diameter (Fig. 4(b)). The piezocone  
 115 tests were performed at a minimum distance of 110 mm from the rigid boundary walls and at a minimum centre  
 116 to centre distance of 75 mm from the anchors (when measured along width  $B$ ). The net cone resistance was  
 117 calculated as  $q_{net} = q_c - \sigma_{vo} + (1 - \alpha_{cone})u_2$ , where  $q_c$  is the measured cone resistance,  $u_2$  is the pore pressure  
 118 measured at the cone shoulder and  $\alpha_{cone}$  is the unequal area ratio calibrated in a pressure chamber as 0.74. A  
 119 velocity of 0.12 mm/s was maintained for these piezocone tests to ensure a drained response, as confirmed from  
 120 the corresponding pore pressure measurements. Correlations of  $R_D$  with  $q_{net}$  (Roy et. al 2019) suggested that  $R_D$   
 121 at the anchor location was approximately ~48% for samples M1 and M2 (similar to the physical measurements)  
 122 and ~65% for sample D1 (within 5% of the physical measurements). The profiles in Fig. 4 indicate that:

- 123 • the spread of  $q_{net}$  across the two medium dense samples (M1 and M2) in Fig. 4 is negligible, suggesting  
 124 relatively similar  $R_D$  values (as shown in Fig. 4(b)) and allowing comparison of tests across these two medium  
 125 dense samples;
- 126 • there was a small increase in  $q_{net}$  over the course of the testing at the anchor embedment depth, by up to 7%  
 127 in the medium dense samples (M1 and M2) and up to 13% in the dense sample (D1); these would correspond  
 128 to a  $R_D$  change of 4.8% and 8.5% in the dense and medium dense samples respectively and are not expected  
 129 to affect the test results significantly.

### 130 *Centrifuge test program*

131 The anchor tests were conducted in the 1.8 m radius beam centrifuge located at the University of Western  
 132 Australia (Randolph et al. 1991) at an acceleration level of 50g. The testing program, as outlined in Table 2,  
 133 involved 26 uplift tests in three methocel saturated sand samples at anchor embedment ratios ( $H/B$ ) in the range  
 134 2.5 to 2.9 and uplift velocities ( $v$ ) varying from 0.01 mm/s to 100 mm/s. The minimum centre to centre distance  
 135 between the plate anchors was 100 mm, established by assuming a pair of slip planes at a very conservative angle  
 136 of 35° (measured to the vertical), thus avoiding overlapping of the failure mechanism.

## 137 **Experimental results**

138 The key experimental results are summarised in Table 2. The mobilised anchor resistance ( $q$ , which is the  
 139 measured load minus the submerged weight of the anchor and rod, divided by the plate area) and corresponding  
 140 pore pressure ( $u$ ) responses for tests at different uplift velocities ( $v$ ) are plotted against normalised displacement  
 141 ( $\delta/B$ ) for all three sand samples in Fig. 5. The peak resistance is also normalised and reported as anchor factors

142 ( $N_\gamma = q_u/\gamma'H$ , where  $\gamma'$  is the submerged unit weight) in subsequent discussions. As evident in Table 2, seven  
143 repeat tests were conducted, with better agreement observed for those in medium dense sand (Samples M1 and  
144 M2 with a maximum variation of 5.6% in peak anchor capacity  $q_u$ ) than those in dense sand (Sample D1 with  
145 maximum variation of 13% in  $q_u$ ), which is consistent with the range in  $q_{net}$  from the cone penetrometer tests.

146 For the slowest tests conducted, with  $v = 0.01$  mm/s in both samples, a purely drained response is obtained,  
147 as the PPT measured hydrostatic pressures. As shown in the inset in Fig. 5(a), there is a rapid increase in  $q$  as  $\delta/B$   
148 increases to 2%, while the drained anchor capacity ( $q_{u,dr}$ ) is subsequently attained within  $\delta/B$  of 15% in medium  
149 dense ( $R_D = 48\%$ ) and 10% in dense sand ( $R_D = 70\%$ ). The drained  $N_\gamma$  values obtained in dense sand are about  
150 18% higher than those for medium dense sand. The  $N_\gamma$  values are compared in Fig. 6 with those obtained from  
151 previous experiments using plate anchors with different widths (or  $B\gamma'/P_a$ , where  $P_a$  is atmospheric pressure) in  
152 the same sand, but in dry conditions (see Roy et al. (2020b)).  $B\gamma'/P_a$  in Fig 6 serves as the stress normalisation  
153 term to account for the different initial stress levels for the tests in saturated and dry sand. The drained  $N_\gamma$  for  $R_D$   
154 = 70% from the current saturated tests match closely with those in dry sand, whereas for the tests with  $R_D = 48\%$ ,  
155 a maximum difference of 10% is observed compared with the tests in dry sand. Also shown in Fig. 6 are the  
156 analytically determined  $N_\gamma$  values for both relative densities, obtained using the limit-equilibrium method in  
157 Giampa et al. (2019). The analytical solution (considered for  $B\gamma'/P_a = 0.1$  with  $\phi$ ,  $\psi_d$  as per Chakraborty & Salgado  
158 2010), provides reasonable agreement with the measurements in medium dense sand ( $R_D \sim 48\%$ ) but overestimates  
159 the measured data by around 28% for the dense sand ( $R_D \sim 70\%$ ). This overestimation is considered to be due to  
160 progressive failure (as explained in Roy et. al (2020b)), that is more dominant in dense sand states than in loose  
161 sand states, and is not accounted for in the analytical solution. As shown later in the paper, better agreement is  
162 obtained when the drained anchor capacity is calculated using the limit equilibrium approach proposed in Roy et  
163 al. (2020b), which encompasses numerically determined anchor capacity factors using the two surface plasticity  
164 model adopted in this paper, and accounts for progressive failure effects.

165 As  $v$  is increased,  $q_u$  increases progressively as the responses change from drained to partially drained and  
166 then towards fully undrained (supported by the pore pressure responses). The corresponding load-displacement  
167 responses show an initial stiff response in  $q$  before a more gradual increase to attain peak resistance, as shown in  
168 the inset in Fig 5. There is remarkable consistency in the gradients of the measured  $q$  responses after the initial  
169 stiff increase (see Fig 5). However, the displacement required to mobilise the peak capacity increases with  
170 increasing  $v$ , and also as the response changes from drained to undrained, with the undrained  $q_u$  mobilised at a

171 displacement that is 3.0 ( $R_D = 70\%$ , D1) and 4.5 - 5.6 ( $R_D = 48\%$ , M1 and M2) times higher than in the  
 172 corresponding drained tests. This may well have implications for offshore design, as the maximum permissible  
 173 displacement criteria may vary from one type of floating structure to another, such that consideration may need  
 174 to be given to both the limiting capacity and the displacement required to mobilise this capacity. The  
 175 corresponding pore pressure responses in both dense and medium dense sample under partially drained and  
 176 undrained conditions show small increases in  $u$  in the initial stages, before strong dilation-induced negative excess  
 177 pore pressure. The shift from positive to negative  $u$  becomes faster with increasing  $R_D$ , with the point of peak  
 178 negative  $\Delta u/u_{hydro}$  coinciding with  $q_u$  ( $\Delta u$  being the measured excess pore pressure).

179 The experimental results (Fig. 5c) also show that the resistance in the medium dense sample ( $R_D = 47\%$ )  
 180 attains a plateau beyond  $v = 60$  mm/s (i.e. test M2\_P100a, M2\_P60a, M2\_T60a in Table 2), indicating that the  
 181 undrained anchor capacity ( $q_{u,un}$ ) has been attained. The corresponding pore pressure responses at the top of the  
 182 plate for these cases also show remarkable consistency, reaching a plateau of  $u \sim -6$  kPa.  $N_\gamma$  values for these tests  
 183 range between 16.55 and 16.9, indicating that the ratio of undrained to drained anchor capacity ( $q_{u,un}/q_{u,dr}$ ) is  
 184 around 5.3. For the tests in dense sand ( $R_D = 70\%$ ), a similar plateau in  $q$  is observed for the tests at  $v = 10$  mm/s  
 185 (D1\_P10a) and  $v = 30$  mm/s (D1\_P30a), with  $q_u$  of 450 kPa and corresponding  $u$  and  $N_\gamma$  values around -40 kPa  
 186 and 17.2 respectively (Fig. 5a). The  $q_u$  values for D1\_P10a and D1\_P30a are 4.6 times higher than the measured  
 187  $q_{u,dr}$  values (i.e. test D1\_P0.01a) at  $R_D = 70\%$ , but only 7% higher than the corresponding  $q_u$  values obtained at a  
 188 uplift velocity,  $v = 10$  mm/s at  $R_D = 48\%$  (i.e. test M1\_P10a, M2\_T10a in Table 2). The  $q_u$  values for D1\_P10a  
 189 and D1\_P30a are 4.6 times higher if compared against the measured  $q_{u,dr}$  values (i.e. with test D1\_P0.01a) in  $R_D$   
 190 = 70%, but only 7% higher if compared against the corresponding  $q_u$  values obtained from  $v = 10$  mm/s for  $R_D =$   
 191 48% (i.e. test M1\_P10a, M2\_T10a in Table 2).

192 The plateau in  $u$  (Fig. 5) suggests the occurrence of cavitation in the undrained anchor tests, as observed  
 193 in triaxial undrained compression tests (McManus & Davis 1997). However, the cavitation pressure,  $u_{cav} \sim -40$   
 194 kPa is higher than the expected theoretical cavitation limit of -100 kPa ( $-1P_a$ ). This may be attributed to the  
 195 complex pore pressure field around the PPT on top of the plate, which were  $0.5B$  away from the plate centre (Fig.  
 196 3). The stress field is clearly dominated by total normal stress towards the centre of the plate, whereas it is  
 197 dominated by shearing along the edges. The consequence is compression leading to positive excess pore pressures  
 198 in the centre (Burns & Mayne 1998, Peuchen et al. 2010), and dilation leading to negative excess pore pressure  
 199 along the edges. The pore fluid may have still cavitated in the shear bands at the edges of the plate, but the average

200  $u$  values measured at the PPT locations remain higher than the theoretical cavitation pressure. Observations of  
201 viscous pore fluid cavitating at a magnitude higher than the theoretical limit of -100 kPa have also been reported  
202 from model tests on gravity platforms on dense sand (Andersen et al. 1994).

203 It is interesting to note that the fastest test, at  $v = 100$  mm/s (D1\_T100a), gave  $q_u = 507$  kPa, which is about  
204 11% higher than those for tests at  $v = 10$  mm/s and 30 mm/s. As pore pressure data were not available for the  
205 fastest test, it is difficult to ascertain whether this increase was a consequence of viscous rate effects (Lehane et  
206 al. 2009, Randolph & Hope 2004, Chow et al. 2020) or a further drop in  $u$ . More experimental data at high  $v$  would  
207 be required to draw an unambiguous conclusion on the precise limits for  $q_{u,un}/q_{u,dr}$  in dense sand.

208 It is important to remember that, although the initial effective stresses were similar in each sample (ignoring  
209 the small changes due to differences in sample density), the initial hydrostatic pressure ( $u_{hydro}$ ) at the soil surface  
210 was lower for the dense sample (D1) than the medium dense sample (M1, M2). This could influence the value of  
211  $q_{u,un}$  for shallowly buried plates, as it will affect the maximum negative pore pressure possible during cavitation  
212 (Andersen et al. 1994). Therefore, the apparently similar range for  $q_{u,un}/q_{u,dr}$  ratios in the two samples (from the  
213 data on Fig. 5) between 4.84 and 5.46 could be attributed to different  $u_{hydro}$  in the two samples. It is worth noting  
214 that  $u_{hydro}$  is also a necessary input (in addition to  $R_D$ ) for calculation of  $q_{u,un}$ , in order to correctly account for the  
215 true dilating potential of the sand.

## 216 **Backbone curve for pure uplift**

217 Consolidation effects on the plate anchor capacity, presented as  $q_u$  against  $V(vd/c_v)$ , are illustrated in Fig.  
218 7(a). The values of  $c_v$  for all samples are listed in Table 2 (from Chow et al. 2019). The drainage path length ( $d$ )  
219 was taken as the radius of an equivalent circle having the same area as the plate anchor (Chung et al. 2006). Fig.  
220 7(a) shows that the drained and undrained capacities were achieved at  $V_{dr} = 1$  and  $V_{un} = 1000$  respectively for both  
221 the medium dense and dense samples. To obtain the backbone curve parameters in Eqn. (1),  $q_u/q_{u,dr}$  ratios are  
222 plotted against  $V$ , which collapses the data from both medium dense and dense sand into a tight band (see Fig.  
223 7(b)). The backbone curve parameters are established (from a regression analysis) as  $q_{u,un}/q_{u,dr} = 5.46$ ,  $V_{50} = 32$   
224 and  $c = 0.71$ .

225 For the horizontal anchors considered here, the value of  $q_{u,un}/q_{u,dr} \sim 5.46$  is much higher than reported  
226 by Chow et al. (2020) for inclined anchors in the same sand ( $q_{u,un}/q_{u,dr} \sim 2.2$ ). This may be explained by the  
227 higher kinematic restraints for the failure mechanism of horizontal anchors, characterised by an inverted

trapezoidal block mechanism (see Fig. A1), compared with bilinear or log-spiral rupture surfaces (Choudhary & Dash 2017, Yue et al. 2020). This is supported by the significantly lower values of  $V_{dr} = 1$  and  $V_{50} = 32$  obtained here, compared with  $V_{dr} = 16$  and  $V_{50} = 175$  for inclined anchors. The rapid rise in mobilised anchor capacity and swift transition into the partially drained response for horizontal anchors suggests that careful consideration of consolidation effects is required in the design of vertically loaded plate anchors, especially for situations where the failure mechanism extends in a quasi-vertical direction to the soil surface.

### Estimation of undrained capacity

The backbone curve for plate anchor uplift (Eqn. (1)) can be effectively utilised if the ratio of the peak anchor capacities (i.e.  $q_{u,un}/q_{u,dr}$ ) is known a priori. The drained capacity (i.e.  $q_{u,dr}$ ) can be estimated using limit equilibrium methods (LEM) such as those in Murray & Geddes (1987) or Vermeer & Sutjiadi (1985), with refinements to these methods proposed by White et al. (2008) and Giampa et al. (2017, 2019) to guide the choice of peak friction and dilation angles, and more recently by Roy et al. (2020b) to better account for the influence of progressive failure. However, estimation of an operative undrained capacity ( $q_{u,un}$ ) at shallow embedment levels remains complicated (Chow et al. 2020) as it depends on the dilation potential of the sand (controlled by  $R_D$  and  $p'$ ) and cavitation. From this standpoint, this section attempts to improve the estimation of  $q_{u,un}$  during cavitation through analytical and finite element (FE) undrained analysis. The FE analysis considers strip anchors for computing efficiency, which can then be extended to other anchor types (e.g. the rectangular anchors in the present centrifuge study) through shape factors.

The relevant analytical expression for computing  $N_{\gamma,un}$  in the case of (dense) dilation induced cavitation scenarios, based on a modified limit equilibrium approach under plane strain conditions (White et al. 2008, Hao et al. 2018) (see Appendix for details on governing equations) can be expressed as:

$$N_{\gamma,un} = N_{\gamma,dr} + \alpha \frac{H}{B} \left( \frac{\Delta u_{max}}{\sigma'_v} \right)^\beta \quad (3a)$$

$$\Delta u_{max} = u_{hydro} + \gamma_w H - u_{cav}, \text{ where } u_{cav} \geq -100 \text{ kPa} \quad (3b)$$

where  $\alpha$  is a soil mobilisation coefficient during uplift and is a function depending on  $R_D$ ,  $\beta$  is a fitting constant,  $u_{hydro}$  is the hydrostatic pressure at soil surface (as defined earlier),  $\gamma_w$  is the unit weight of water and  $\Delta u_{max}/\sigma'_v$  is a normalised maximum pore pressure change according to the limiting pore pressure before cavitation sets in. Eqn. (3a) and (3b) are valid provided that the absolute magnitude of  $\Delta u_{max}$  remains below the limit at which the

dilation induced, negative excess pore pressure result in critical state (constant volume) conditions. A theoretical estimate of this ‘critical state’  $\Delta u_{max}$  can be readily obtained using the concepts provided in Fear & Robertson (1995). With the current CSL parameters for UWA sand and assuming even a modest  $R_D = 40\%$  at  $H/B = 3$ , a back analysis shows that this limit is not exceeded at least for 120 metres of water depth and remains valid for even higher water depths if  $R_D$  is higher. The term  $\Delta u_{max}/\sigma'_v$ , introduced to account specifically for different water depths through  $u_{hydro}$  (representing water depth), governs the relative magnitude of  $N_{\gamma,un}$  with respect to  $N_{\gamma,dr}$ . It is important to observe that  $u_{hydro}$  would be a low value in shallow waters, such that the term  $\Delta u_{max}/\sigma'_v$  ( $\sigma'_v$  being equal to  $\gamma'H$ , which has almost similar magnitude to  $\gamma_w H$ ) would attain a value close to 1.0 in shallow water and increase progressively for deeper water at a given  $\sigma'_v$ . The operative values of  $\alpha$  and  $\beta$  in Eqn. (3) are obtained through undrained FE analysis of strip anchors, as explained in the next section and summarised in Eqn. (31).

### *Numerical model details*

The FE analysis was undertaken in Abaqus (Dassault Systemes, 2012) for strip anchors in sand (plane strain) with an unstructured mesh of linear quadrilateral elements (CPE4) (Fig. 8). The FE results were validated with the experiments involving rectangular anchors by applying appropriate shape factors ( $S_f$ ). The analysis considered a strip anchor with  $B = 1$  m, varying  $H/B$  between 2 and 4,  $R_D$  between 30% and 85%, with initial hydrostatic pressure ( $u_{hydro}$ ) at the soil surface between 0 and 50 kPa. Bottom and lateral boundaries of the domain from the centre of the plate were varied in the range of  $10B$  to  $12B$  and  $8B$  to  $10B$  respectively, depending on  $H/B$ . The anchor, with slightly filleted corners, was defined as weightless and fully rigid, with its centre taken as the reference point for the application of displacements. The interface between the plate and the soil was modelled as fully rough. The coefficient of earth pressure  $K_0$  was taken as 0.5. An effective stress analysis was achieved in Abaqus using the equivalent total stress analysis approach as adopted in Yi et al. (2012) and Chang (2018), where the pore water stiffness from the bulk modulus of water ( $K_w$ ) was superimposed into the incremental effective stress stiffness matrix to compute the resulting  $\Delta u$  changes. The changes in  $\Delta u$  at each of the Gauss points was stored as a state variable during the process.

Initial trials showed that the magnitude of  $u$  in shear bands and in the zones below the plate can attain large unrealistic negative values lower than the theoretical cavitation limit of water. Once cavitation sets in, the water should change to vapour, limiting further reduction in pore water pressure and releasing the constraint on constant volume (Gawin & Sanavia 2010, Mcmanus & Davis 1997). Hence, a pore pressure cavitation cut-off was introduced to allow realistic estimates of capacity. This was implemented numerically by tracking  $u$  at the Gauss

284 points through a state variable, and by allowing a smooth transition of  $K_w$  from its true value of 2 GPa (for  
285 undrained conditions) to 0 (fully drained conditions) through a power law reduction of the form  $\langle 1 - \left( \frac{\langle -u \rangle}{\langle -u_{cav} \rangle} \right)^{30} \rangle$   
286 based on the current value of  $u$ , where  $\langle \rangle$  are Macaulay brackets, such that  $\langle x \rangle = x$  if  $x > 0$  and  $\langle x \rangle = 0$  if  $x < 0$ .  
287 Such an expression keeps  $K_w$  unaltered for positive pore pressure scenarios ( $u > 0$ ) and prevents further build-up  
288 of  $u$  once  $u_{cav}$  is attained.

### 289 *Constitutive model details*

290 The soil constitutive model adopted in the FE simulations is the two surface plasticity model proposed by  
291 Dafalias & Manzari (2004) or Dafalias et al. (2004) as it embraces typical bounding surface plasticity features  
292 with kinematic hardening. The model controls sand behaviour with simultaneous interaction of the yield surface  
293 with three other surfaces, namely the critical state, bounding and dilatancy surfaces, thus allowing realistic  
294 predictions of stresses and strains over a wide range of densities and  $p'$  levels. A unique CSL in  $e - p'$  space (Li &  
295 Dafalias 2012) is used in the current model. Some minor practical modifications, pertaining to the treatment of  
296 fabric anisotropy as explained in Roy et al. (2020a), were introduced through scaling factors ( $f$  and  $f_h$ , see Eqn.  
297 (27) to (29) in Table 3) to adjust the influence of the state parameter ( $\psi$ ) on the model governing equations.  
298 Horizontally isotropic sand was assumed with appropriate inputs for the fabric tensor as per Eq. (24) (and  $\mathbf{F}$  as in  
299 Eq. (25)). The complete details, calibration procedure against triaxial and simple shear tests and subsequent FE  
300 implementation of the model can be found in Roy et al. (2020a) and a brief summary of the model parameters is  
301 provided in Table 3.

### 302 *Validation results and load deformation plots*

303 The FE analysis is validated in Fig. 9 using the centrifuge tests conducted at  $H/B = 2.7$ ,  $R_D = 50\%$  and  $70\%$   
304 with the surface values of  $u_{hydro}$  being taken as 20 kPa and 10 kPa respectively (as per the experiments). A number  
305 of trial simulations indicate that cavitation of the pore fluid below the plate occurs at displacement levels of  $0.04B$   
306 (i.e. the displacement level associated with mobilisation of peak undrained capacity). On this basis, the current  
307 simulations allow separation between the plate and the soil underneath the anchor, such that the numerically  
308 determined  $N_{\gamma,un}$  values reflect a condition where the response at the underside of the plate anchor is drained.  
309 Using the full theoretical cavitation pressure ( $u_{cav}$ ) in the simulations of -100 kPa gave an anchor capacity that  
310 was on an average 60% higher than the experimental data. From a design perspective, the value of -50 kPa adopted  
311 in the simulations appears reasonable (and conservative) but requires further investigation in future work.

312 The simulated peak values of  $N_{\gamma,un}$  was 11.88 at  $\delta/B = 5\%$  for  $R_D = 50\%$  (with  $u_{hydro} = 20$  kPa) and  $N_{\gamma,un} =$   
 1 11.22 for  $R_D = 70\%$  at  $\delta/B = 3.6\%$  (with  $u_{hydro} = 10$  kPa) respectively. Similar to the trend observed in the  
 2 313 experiments, the corresponding undrained limiting capacity is attained at  $\delta/B$  levels which are 2.0 ( $R_D = 70\%$ ) to  
 3 314 3.3 ( $R_D = 50\%$ ) times higher than the corresponding drained simulations. However, the magnitude of the mobilised  
 4 315 displacements in the experiments are approximately ~10-12 times higher than in the simulations, independent of  
 5 316 the drainage response. This is partly due to particle scale effects along the shear bands (Ng and Springman 1994,  
 6 317 Palmer et al. 2003), but also because the current model uses a lower operative elastic shear modulus than the true  
 7 318  $G_{max}$  of the soil and does not consider degradation of elastic shear modulus with strain level (Salgado et al. 2000).  
 8 319 Although better agreement with the measured stiffness can be obtained by adjusting the stiffness parameters, this  
 9 320 would worsen the agreement with the anchor capacity as adjusting  $G_0$  also affects the plastic response in the model  
 10 321 (Roy et al. 2020a).  
 11 322

23 323 To enable comparison with the experimental results, the shape factor ( $S_f$ ) proposed by Dickin (1988) was  
 24 324 used to scale the measured results for the rectangular anchors to those for the strip anchor, according to

$$27 \quad 28 \quad 29 \quad 30 \quad 31 \quad 32 \quad 33 \quad 34 \quad 35 \quad 36 \quad 37 \quad 38 \quad 39 \quad 40 \quad 41 \quad 42 \quad 43 \quad 44 \quad 45 \quad 46 \quad 47 \quad 48 \quad 49 \quad 50 \quad 51 \quad 52 \quad 53 \quad 54 \quad 55 \quad 56 \quad 57 \quad 58 \quad 59 \quad 60 \quad 61 \quad 62 \quad 63 \quad 64 \quad 65$$

$$325 \quad S_f = 1 - \frac{jB}{3LH}(6B - 7H) \quad (30a)$$

30 326 where  $j$  was found by Dickin (1988) to be varying from 0.4 at  $R_D = 33\%$  to 1.0 at  $R_D = 76\%$  for anchor uplift tests  
 31 327 in Erith sand. As the current experiments involved slightly different relative densities, a more representative  $S_f$   
 32 328 can be achieved by linearly interpolating between the limits 0.4 and 1.0 for  $R_D$  of 33% and 76% as per Eqn. (30b):

$$37 \quad 38 \quad 39 \quad 40 \quad 41 \quad 42 \quad 43 \quad 44 \quad 45 \quad 46 \quad 47 \quad 48 \quad 49 \quad 50 \quad 51 \quad 52 \quad 53 \quad 54 \quad 55 \quad 56 \quad 57 \quad 58 \quad 59 \quad 60 \quad 61 \quad 62 \quad 63 \quad 64 \quad 65$$

$$329 \quad j = 0.0132 R_D(\%) - 0.013 \quad (30b)$$

40 330 For  $R_D \sim 50\%$  and  $\sim 70\%$  in the current experiments, Eqn. (30b) gives  $j = 0.65$  and  $0.91$  respectively, and ultimately  
 41 331 results in  $S_f$  values of 1.52 for medium dense sand and 1.73 in dense sand as per Eqn. (30a). These factors result  
 42 332 in equivalent strip anchor  $N_{\gamma,un}$  values from the experiments of 11.12 (= 16.9/1.52) and 10.65 (= 18.42/1.73) for  
 43 333 the tests at  $R_D = 50\%$  and  $70\%$  respectively; these are 6.8% and 5.35% lower than the FE simulated values. In  
 44 334 fact, if an  $S_f$  of 1.63 at  $H/B = 2.7$  had been chosen for the rectangular plates as per the studies in Frydman &  
 45 335 Shaham (1989) of buried slab/anchors, the resulting capacity estimates for  $R_D = 50\%$  and  $70\%$  (considering an  $S_f$   
 46 336 = 1.63) would also have been within margins of 12.6% and 9.4% for the two cases. The improved comparison  
 47 337 justifies the use of an operative  $u_{cav}$  of -50 kPa for these simulations.

57 338 Fig. 10 compares the respective failure mechanisms at peak capacity in  $R_D = 70\%$  at  $H/B = 2.7$  between  
 58 339 the undrained (with cavitation) and drained uplift cases in terms of the normalised soil displacements ( $\delta_{soil}/\delta$ )

340 contours and normalised octahedral plastic shear strain contours ( $\Delta\gamma_{oct}^p/\Delta\gamma_{oct,max}^p$ ). The displacement contours in  
 341 Fig. 10(a) shows that the influence zone mobilised at peak capacity for the undrained case is larger than for drained  
 342 conditions, with the corresponding width of the influence zone at the soil surface for undrained uplift being 20%  
 343 higher than that for drained uplift. On the other hand, Fig. 10(b), comparing the corresponding normalised  
 344 octahedral plastic shear strain contours ( $\Delta\gamma_{oct}^p/\Delta\gamma_{oct,max}^p$ ) mobilised at peak capacity in  $R_D = 70\%$ , shows similar  
 345 inclination of shear bands of  $20^\circ \pm 3^\circ$  and  $18^\circ \pm 4^\circ$  for the undrained and drained cases respectively, which is  
 346 slightly higher than a value of  $15.6^\circ$  that would have been obtained using Bolton's (1986) correlations at respective  
 347 in-situ stress levels. The minor differences in the inclination of the shear bands justifies the use of the same dilation  
 348 angle in the analytical expression for the drained and undrained cases (as elaborated in the Appendix).

349 Fig. 11 demonstrates the proximity of the pore fluid to cavitation at peak  $N_\gamma$  for  $H/B = 2.7$  in  $R_D = 70\%$   
 350 with  $u_{hydro}$  of 10 kPa. The figure plots the excess pore pressure ( $\Delta u$ ) normalised with respect to the maximum  
 351 (absolute) change possible at that location prior to cavitation ( $\Delta u_{cav}$ ). Distinct bands of cavitating pore fluid zones  
 352 (values close to -1.0) develop along the edges of the plate, mostly coinciding with the shear bands discussed above.  
 353 There is also a zone of positive excess pore pressure on top of the plate with values higher than 1. This explains  
 354 why the spot measurements of pore pressure in the experiments were lower than the theoretical value of  $u_{cav}$ . Fig  
 355 11(b) plots  $p'$ ,  $q$  and  $\Delta u$  for two different points located a vertical distance  $0.4B$  above the plate:  $X_1$  which is  
 356 directly above the centre of the plate and  $X_2$  which is located in the zone of intense shearing. With increasing  $\delta/B$ ,  
 357 the element at  $X_1$  clearly shows evidence of minor dilation, whereas cavitation is evident at  $X_2$  with  $\Delta u$  attaining  
 358 constancy around -80 kPa (i.e.  $\Delta u_{cav}$  at that location). The corresponding  $p'$  and  $q$  increase sharply during the initial  
 359 phase of fully undrained loading, but drop in magnitude following the initiation of cavitation, with a switch to a  
 360 drained (simple shear) loading (i.e.  $K_w$  attaining 0).

#### 361 *Numerical results from the parametric study*

362 The above numerical approach has then been extended parametrically to compute  $N_{\gamma,un}$  for  $R_D$  between  
 363 30% and 85%,  $H/B$  between 2 and 4 and normalised pore pressure variation potential ( $\Delta u_{max}/\sigma'_v$ ) between 2 and  
 364 8.7. For undrained uplift conditions, a change of scale (i.e. changing  $B$  but with the same  $H/B$  etc.) would result  
 365 in a proportionate change in the values of  $\Delta u_{max}/\sigma'_v$ , such that the numerical investigation does not need repeating  
 366 over different stress levels. Fig. 12(a) plots values of  $N_{\gamma,un}$  against different  $\Delta u_{max}/\sigma'_v$ , showing that  $\Delta u_{max}/\sigma'_v$  has  
 367 a profound influence on  $N_{\gamma,un}$ . In fact, at values of  $\Delta u_{max}/\sigma'_v$  close to 2 (expected in very shallow water as explained  
 368 earlier),  $\Delta u_{max}/\sigma'_v$  seems to be the only influencing factor for  $N_{\gamma,un}$  with  $R_D$  and  $H/B$  having inconsequential effects.

369 This observation is consistent with the centrifuge test results that suggest similar  $N_{\gamma,un}$  numbers for  $R_D = 50\%$  and  
 1  
 2 370  $70\%$  at  $\Delta u_{max}/\sigma'_v$  levels of around 3.5. This suggests that, if a plate anchor is designed to operate in conditions  
 3  
 4 371 having periodic variations in water level, extra care needs to be taken in the design.  
 5

6  
 7 372 The results of the FE parametric study can now be used to determine  $\alpha$  and  $\beta$  in Eqn. (3). This is achieved  
 8  
 9 373 by plotting  $(N_{\gamma,un} - N_{\gamma,dr})/(H/B)$  (referred to here as the depth normalised capacity difference) against  
 10  
 11 374  $\Delta u_{max}/\sigma'_v$  as shown in Fig. 12(b). Fig 12b shows that there is a linear relationship between depth normalised  
 12  
 13 375 capacity difference and  $\Delta u_{max}/\sigma'_v$  for different values of relative density, with higher dependence for higher values  
 14  
 15 376 of  $R_D$  reflected in the higher  $\alpha$  values, which indirectly influences  $N_{\gamma,un} \cdot N_{\gamma,un}$ . The effect becomes appreciable  
 16  
 17 377 for  $\Delta u_{max}/\sigma'_v$  greater than about 5 (i.e. in deep water relative to the anchor embedment), beyond which the range  
 18  
 19 378 for  $N_{\gamma,un} - N_{\gamma,dr}$  exceeds 20%. The values of peak drained anchor factors (i.e.  $N_{\gamma,dr}$ ) have been obtained using  
 20  
 21 379 the modified LEM approach based on the same numerical model as described in Roy et al. (2020b), that directly  
 22  
 23 380 accounts for stress level dependency and progressive failure. The values of  $\alpha$  and  $\beta$  for Eqn. (3) were fitted through  
 24  
 25 381 least square regression to give:

$$382 \quad \alpha = (0.01 * R_D(\%))^0.2 \quad (31a),$$

$$383 \quad \beta = 1.11 \quad (31b)$$

384 Now Eqn. (3), in conjunction with Eqn. (31) and the backbone curve (Eqn. (1)) can be used to predict anchor  
 385 capacity for different drainage regimes in sand. Further experiments in sands of varying dilative characteristics  
 386 would assist in quantifying the potential variation of the constants  $\alpha$  and  $\beta$  used in Eqn. (3).  
 387

388 The versatility of the backbone curve (Eqn. (1)) and the analytical approach for undrained capacity (Eqn.  
 389 (30) is evaluated in Fig. 13 by comparing predicted  $N_\gamma$  values against measured values in medium dense ( $R_D =$   
 390  $50\%$ ) and dense ( $R_D = 70\%$ ) sand. Fig. 13 includes the entire set of measured  $N_\gamma$  values from the 26 uplift tests  
 391 reported in Table 2, reduced to their equivalent strip anchor values by using  $S_f$  as described in Dickin (1988)  
 392 ( $\sim 1.52 - 1.73$ ) or Frydman & Shaham (1989) ( $\sim 1.63$ ). For use of Eqn. (1), the predicted  $N_{\gamma,dr}$  values have been  
 393 computed using Eqns. (A1-A4) in the Appendix (further details are provided in Roy et al. (2020b)).  $N_{\gamma,un}$  has  
 394 been obtained from Eqn. (3) using  $u_{hydro} = 10$  kPa and 20 kPa for  $R_D = 70\%$  and  $R_D = 50\%$  respectively, and  $u_{cav}$   
 395  $= -50$  kPa (as assumed for comparisons with the experiments). Using this approach, the predicted values of  
 396  $N_{\gamma,un}$  turn out to be 11.55 and 12.5 for  $R_D = 70\%$  and  $R_D = 50\%$  respectively, compared to the shape adjusted  
 measured values (i.e.  $N_{\gamma,un}/S_f$ ) which are in the range 10.18 - 11.21 for  $R_D = 70\%$  and 9.36 - 10.35 and  $R_D = 50\%$ .

397 The predicted ratios of  $q_{u,un}/q_{u,dr}$  turn out to be 4.34 and 4.84 for  $R_D = 70\%$  and  $R_D = 50\%$  respectively, which  
398 are about 6% and 13% lower than the (strip converted) experimental values. These ratios of  $q_{u,un}/q_{u,dr}$  then allow  
399 subsequent prediction of  $N_\gamma$  for any given  $V$  through Eqn. (1).

## 400 **Conclusions**

401 This study has investigated the effect of displacement rate, and hence consolidation, on horizontal plate  
402 anchor capacity in medium dense and dense sand through centrifuge experiments covering four orders of  
403 magnitude of variation in dimensionless velocity ( $V = 0.2$  to 1930). As  $V$  increases in the experiments, the  
404 mobilised anchor capacity increases progressively, with pore pressure responses on top of the plate becoming  
405 increasingly negative (i.e. suction). At  $V > 1000$ , the anchor capacity and pore pressure are observed to reach a  
406 plateau, indicating undrained conditions with occurrence of cavitation. However, the measured pore pressures  
407 remain higher than the theoretical  $u_{cav}$  of -100 kPa, due to the complex pore pressure field at the PPT location.  
408 The consolidation effects on anchor capacity can be captured using a backbone curve giving  $q_{u,un}/q_{u,dr}$  as a function  
409 of  $V$  (Eqn. (1)) with fitted parameters of  $q_{u,un}/q_{u,dr} = 5.46$ ,  $V_{50} = 32$  and  $c = 0.71$ .

410 For application of the backbone curve in practice, estimation of the undrained anchor capacity ( $q_{u,un}$ )  
411 remains the biggest uncertainty due to the pore pressure response, particularly upon cavitation. In order to reduce  
412 these uncertainties in estimating  $q_{u,un}$ , a finite element parametric study was undertaken using a bounding surface  
413 plasticity model, taking into account the effects of sand relative density, hydrostatic pressure (controlling  
414 cavitation potential) and embedment ratio. The FE study showed that the magnitude of the initial hydrostatic  
415 pressure has a profound influence on undrained anchor capacity, meaning that particular care should be exercised  
416 when designing plate anchors for use in shallow water.

417 The influence of the governing factors in undrained plate anchor uplift, as obtained in the parametric study,  
418 was represented systematically through a simple analytical model (Eqn. (3)). With proper choice of  $u_{cav}$   
419 (influencing the contribution of  $\Delta u_{max}/\sigma'_v$ ), the combined numerical and analytical solution was shown to agree  
420 reasonably well with the experimental results. Further work is needed to understand the deviation of  $u_{cav}$  from the  
421 theoretical value, and to further validate the proposed model on a wider range of sands including natural deposits.

## 422 Appendix

1  
2  
3 423 This section derives an analytical expression for the undrained capacity ( $N_{\gamma,un}$ ) of a horizontal strip anchor  
4  
5 424 in cavitating sand. The approach follows the limit equilibrium approach for drained capacity ( $N_{\gamma,dr}$ ) assuming an  
6  
7 425 inclined slip model (White et al. 2008, Tom et al. 2017), but with modifications to account for the additional  
8  
9 426 contribution of excess pore pressure. For the inverted trapezoidal mechanism shown in Fig A1, the peak drained  
10  
11 427 uplift resistance force (referred here as  $P_{dr}$ ) is approximated as the weight of the rigid (trapezoidal) block above  
12  
13 428 the anchor and the shear resistance on the slip planes. Based on Fig A1,

$$16 \quad 429 \quad P_{dr} = HB\gamma' + \gamma'H^2 \tan(\psi_d) + 2 \int_0^H (\tau_f \cos \psi_d - \sigma'_n \sin \psi_d) \frac{dz}{\cos \psi_d} \quad (A1)$$

19  
20 430 Expressing  $\tau_f = K_1 \sigma'_n$  and  $\sigma'_n = K_2 \sigma'_v$ , Eqn. (A1) becomes

$$23 \quad 431 \quad P_{dr} = HB\gamma' + \gamma'H^2 [\tan(\psi_d) + K_1 K_2 - K_2 \tan(\psi_d)] \quad (A2)$$

$$26 \quad 432 \quad \text{Finally, } N_{\gamma,dr} = \frac{P_{dr}}{HB\gamma'} = 1 + \frac{H}{B} F_{u,ps,dr} \quad (A3)$$

29 433 where  $F_{u,ps,dr} = [\tan(\psi_d) + K_1 K_2 - K_2 \tan(\psi_d)]$ . A series of drained numerical analyses in Roy et al. (2020b)  
30  
31 434 showed that the drained uplift factor  $F_{u,ps,dr}$  can be represented as:

$$34 \quad 435 \quad F_{u,ps,dr} = 0.43 + 0.052I_R \quad (A4)$$

37 436 where  $I_R = (= \frac{R_D(\%)}{100} (10 - \ln p') - 1)$  is a relative dilatancy index computed at a mean stress level of  
38  
39 437  $p' = \sigma'_{vo}(0.25 H/B + 0.5)$ .

42 438 For undrained conditions, the shear resistance on the slip planes is modified to account for the additional  
43  
44 439 contribution of the excess pore pressures ( $\Delta u$ ), resulting in  $\sigma'_n = K_2(\sigma'_v - \Delta u)$ . Due to cavitation of the pore  
45  
46 440 fluid beneath the plate, potential tensile resistance below the plate is not included, such that the calculated  
47  
48 441 undrained capacity is conservative and also serves to make the analytical model an equivalent of the numerical  
49  
50 442 model reported in the ‘Numerical model details’ section. Therefore, the peak undrained uplift resistance force  
51  
52 443 (referred to here as  $P_{un}$ ) may be calculated using a modified form of Eqn. (A2):

$$55 \quad 444 \quad P_{un} = BH\gamma' + \gamma'H^2 [\tan(\psi_d) + K_1 K_2 - K_2 \tan(\psi_d)] - 2H (\Delta u)(K_1 K_2 - K_2 \tan(\psi_d)) \quad (A5)$$

58 445 such that the undrained capacity factor ( $N_{\gamma,un}$ ) may be written as:

$$N_{\gamma,un} = \frac{P_{un}}{HB\gamma'} = N_{\gamma,dr} - 2 \frac{H(\Delta u)}{B\sigma'_v} (K_1 K_2 - K_2 \tan(\psi_d)) \quad (A6)$$

In case of dilation induced cavitation, a term  $\Delta u_{max}$  is introduced (referred to as the pore pressure variation potential), that indicates the maximum negative pore pressure possible at the anchor location before the pore fluid cavitates.  $\Delta u_{max}$  is influenced by the value of the initial hydrostatic pressure ( $u_{hydro}$ ) at the soil surface and cavitation pressure ( $u_{cav}$ ). The value of  $\Delta u_{max}$  governs the maximum change in effective stress possible at the anchor location under fully undrained conditions. Therefore,

$$\Delta u_{max} = u_{hydro} + \gamma_w H - u_{cav} \quad (A7)$$

Eqn. (A6) can then be written as:

$$N_{\gamma,un} = N_{\gamma,dr} + 2 \frac{H}{B} \alpha \frac{\Delta u_{max}}{\sigma'_v} (K_1 K_2 - K_2 \tan(\psi_d)) \quad (A8)$$

where  $\Delta u/\sigma'_v = -\alpha \Delta u_{max}/\sigma'_v$  (noting that the minus sign is required as  $\Delta u$  will be negative during dilation) and  $\Delta u_{max}$  is calculated according to Eqn. (A7). The term  $(K_1 K_2 - K_2 \tan(\psi_d))$  in Eqn. (A8) is density and stress level dependent. The additional density dependence can be introduced by introducing a  $R_D$  dependence on the term  $\alpha$ , whereas the flexibility of additional stress level dependence may be introduced by incorporating a constant  $\beta$  on the stress level term (controlled by  $\Delta u_{max}/\sigma'_v$ ).  $\alpha$  therefore serves as a soil mobilisation coefficient during uplift, whereas  $\beta$  serves as fitting constant. Therefore, Eqn. (A8) can be rewritten in the form:

$$N_{\gamma,un} = N_{\gamma,dr} + \alpha \frac{H}{B} \left( \frac{\Delta u_{max}}{\sigma'_v} \right)^\beta \quad (A9)$$

where  $\alpha = f(R_D)$  and  $\beta$  is a fitting constant. Now, as both the analytical model and the FE model encompass the same geotechnical phenomena, the constants  $\alpha$  and  $\beta$  can be obtained by best-fitting them over the numerical determined previously  $N_{\gamma,un}$  from FE analyses.

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53 **Table captions**

56 625 Table 1: Physical properties of UWA Silica sand (Chow et al. 2019)

58  
59 626 Table 2: Centrifuge testing program

627 Table 3: Model details

628 Table 4 : Model parameters for UWA silica sand

## 629 **Figure captions**

630 Fig. 1 : Backbone curve framework showing the variation in resistance across different consolidation  
631 regimes

632 Fig. 2 : Variation of normalised viscosity with temperature for methocel F450

633 Fig. 3 : Centrifuge test details: (a) test setup (not to scale), (b) plate anchor model, (c) “wished in place”  
634 anchors during sample preparation (d) piezocone testing in saturated sand for sample characterisation

635 Fig. 4 : Profiles of : (a) net cone tip resistance and (b) relative density estimated using the empirical  
636 correlations in Roy et al. (2019)

637 Fig. 5 : Mobilised capacity and pore pressure response at different uplift velocities in (a) dense sand of  $R_D$   
638 = 70% (b) medium dense sand of  $R_D = 48%$  (c) medium dense sand of  $R_D = 47%$

639 Fig. 6 : Comparison of measured  $N_\gamma$  for rectangular plates in dry and saturated drained conditions for (a)  
640 dense sand of  $R_D = 70\% - 76%$  (b) medium dense sand of  $R_D = 40\% - 50%$

641 Fig. 7 : Changing anchor capacity with dimensionless velocity: (a) mobilised anchor capacity ( $q_u$ ) and (b)  
642 normalised mobilised anchor capacity ( $q_u/q_{u,dr}$ )

643 Fig. 8: Typical mesh details and boundary conditions for the anchor problem in Abaqus ( $H/B = 2.7$ )

644 Fig. 9: Comparison of simulated and measured anchor factors (adjusted by shape factors) at  $H/B = 2.7$  in:  
645 (a) dense sand ( $R_D = 70\%$ ) (b) medium dense sand ( $R_D = 50\%$ )

646 Fig. 10: Comparison of numerically determined: (a) failure mechanisms, (b) shear strain distributions at  
647  $H/B = 2.7$  in dense sand ( $R_D = 70\%$ ) for undrained and drained cases

648 Fig. 11: (a) Contours of excess pore pressure ( $\Delta u$ ) normalised by the maximum possible pore pressure  
649 change possible prior to cavitation ( $\Delta u_{cav}$ ) at  $H/B = 2.7$  in  $R_D = 70\%$  with  $u_{hydro} = 10$  kPa; (b) variation of  $p'$ ,  $q$  and  
650  $\Delta u$  at two different points in the domain

651 Fig. 12: Variation of: (a)  $N_\gamma, u_n$  and (b) depth normalised relative difference,  $(N_\gamma, u_n - N_\gamma, dr)/(\frac{H}{B})$  from  
652 the FE simulations in dense and medium dense sands

653 Fig. 13: Comparison of measured and predicted anchor factors with changing dimensionless velocity in

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4 655 Fig. A1: Schematic representation of: (a) plate anchor uplift at failure (b) failure mechanism with inclined

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Table 1: Physical properties of UWA Silica sand (Chow et al. 2019)

Soil properties	UWA Silica sand
Angularity	Sub-rounded to sub-angular
Maximum void ratio( $e_{max}$ ) <sup>1</sup>	0.789
Minimum void ratio( $e_{min}$ ) <sup>1</sup>	0.512
Specific gravity <sup>2</sup>	2.67
Uniformity Coefficient ( $U$ )	1.73
Particle mean size ( $d_{50}$ )	0.2 mm

<sup>1</sup> Maximum and minimum void ratios calculated as per AS 1289.5.5.1

<sup>2</sup> Specific gravity of soil particles calculated as per AS 1289.3.5.1

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Table 2: Centrifuge testing program

Sample number and $R_D$	Unit weight, $\gamma'$ (kN/m <sup>3</sup> )	Test ID*	$H/B$	Velocity, $v$ (mm/s)	Temperature, $T$ (°C)	$V$	Peak mobilised anchor capacity, $q_u$ (kPa)	Peak $N_\gamma = q_u/\gamma H$
D1 ( $c_v \sim 4.3 \times 10^{-4}$ m <sup>2</sup> /s, $R_D = 70\%$ )	10.25	D1_P0.01a	2.6	0.01	-†	0.2	97.27	3.80
		D1_T0.1a	2.8	0.1	19.2	1.9	174.46	6.34
		D1_T0.3a	2.65	0.3	18.7	6.3	211.65	8.28
		D1_P1a	2.55	1	20.6	20.2	215.58	8.59
		D1_T3a	2.7	3	20.4	57.8	365.78	13.82
		D1_T3b	2.5	3	19.2	65.4	317.73	12.92
		D1_P10a#	2.6	10	21.1	194.2	446.97	17.46
		D1_P30a	2.63	30	20.7	585.2	450.98	17.47
D1_T100a	2.8	100	19.4	1933.7	506.81	18.42		
M1 ( $c_v \sim 7 \times 10^{-4}$ m <sup>2</sup> /s, $R_D = 48\%$ )	9.87	M1_P0.01a	2.94	0.01	20.7	0.1	93.76	3.36
		M1_T0.01a	2.88	0.01	20.5	0.1	103.25	3.74
		M1_P0.1a	2.93	0.1	20.8	1.2	120.47	4.34
		M1_T1a	2.95	1	21.1	11.5	222.06	7.95
		M1_T3a	2.84	3	20.8	36.0	272.18	10.10
		M1_P10a	2.88	10	20.9	118.2	363.9	13.34
		M1_T100a	2.98	100	21.0	1139.4	467.43	16.55
M2 ( $c_v \sim 6.4 \times 10^{-4}$ m <sup>2</sup> /s, $R_D = 47\%$ )	9.86	M2_P0.01a	2.63	0.01	21	0.1	79.1	3.17
		M2_T0.01a	2.66	0.01	21.3	0.1	83.8	3.32
		M2_P0.1a	2.6	0.1	21.2	1.3	113.2	4.59
		M2_P1a	2.6	1	21.3	12.9	169.7	6.89
		M2_T1a	2.8	1	21.3	12.0	193.5	7.3
		M2_P10a	2.6	10	21.3	125.3	317.3	12.88
		M2_T10a	2.67	10	21.3	129.1	343.8	13.56
		M2_P60a	2.7	60	21.3	744.6	432.0	16.9
M2_T60a	2.61	60	21.34	768.3	404.4	16.31		
M2_P100a	2.65	100	21.2	1270.4	420.6	16.75		

\* The tests are denoted 'Ax1\_Wx2\_m', where 'A' is either 'D' or 'M' for a dense or medium dense sample respectively, x1 is the sample number, 'W' is either 'P' or 'T' to denote plates with and without a PPT respectively, x2 is the uplift velocity in mm/s and 'm' denotes the test recurrence ('a' for the first test, 'b' for the second etc.).

† Data not available.

# LVT slip during test

Table 3: Model details

Parameters	Description	Governing equations (Eqn. No)	Notes/Reference
$e_o$	Controls the intercept of the CSL at very low mean stress levels in $e - (p'/P_a)^\xi$ space	$e_c = e_o - \lambda_c \left( \frac{p'}{P_a} \right)^\xi \quad (4)$ $\psi = (e - e_c) \quad (5)$	Li & Wang (1998), Dafalias & Manzari (2004)
$\lambda_c$	Controls the slope of the CSL in $e - (p'/P_a)^\xi$ space		
$\xi$	Influences the linearisation of the CSL with $p'$		
$M$	Controls the critical state stress ratios in triaxial compression		
$c_l$	Controls the critical state stress ratios in triaxial extension		
$G_o$	Elastic shear modulus	$G = G_o P_a \frac{(2.97-e)^2}{(1+e)} \left( \frac{p'}{P_a} \right)^{1/2} \quad (6)$	Dafalias & Manzari (2004)
$\nu$	Poisson's ratio	$K = G \frac{2(1+\nu)}{3(1-2\nu)} \quad (7)$ $d\varepsilon_e^e = ds/2G \quad (8)$ $d\varepsilon_v^e = dp/K \quad (9)$	
$m$	Controls the size of the yield surface	$f = [(\mathbf{r} - \boldsymbol{\alpha}) : (\mathbf{r} - \boldsymbol{\alpha})]^{1/2} - \sqrt{2/3} m = 0 \quad (10)$	
$n_d, n_b$	Influences the mobilised stress ratio at peak friction angle and at phase transformation	$\boldsymbol{\alpha}_\theta^a = \sqrt{\frac{2}{3}} (M_\theta^a - m) \mathbf{n} \quad (a = b, d, c) \quad (11)$ $M_\theta^c = g_1(\theta, c_1) M \quad (12)$ $g_1(\theta, c_1) = \left( \frac{2c_1^4}{((1+c_1^4) - (1-c_1^4) \cos 3\theta)} \right)^{1/4} \quad (13)$ $M_\theta^b = [M_\theta^c \exp(-n^b \psi_f)] \quad (14)^*$ $M_\theta^d = [M_\theta^c \exp(n^d \psi_f)] \quad (15)^*$	Dafalias & Manzari (2004), Sheng et al. (2000)

1 2 3 4 5 6 7 8 9	$c_h, h_{in}$	Plastic modulus constants	$K_p = \frac{2}{3} p' h (\alpha_0^b - \alpha) : \mathbf{n} \quad (16)$ $h = \frac{b_o}{(\alpha - \alpha_{in}) : \mathbf{n}} \quad (17)$ $b_o = h_o G_o (1 - c_h e) \left( \frac{p'}{P_a} \right)^{-1/2} \quad (18)$ $h_o = h_{in} f_h \quad (19)^*$	Dafalias et al. (2004)
10 11 12 13 14 15 16 17	$A_o$	Parameter affecting the rate of dilatancy	$d\varepsilon^p = \langle L \rangle \mathbf{R} \quad (20)$ $\frac{dg}{d\sigma} = \mathbf{R} = \mathbf{R}' + \frac{1}{3} D \mathbf{I} = B \mathbf{n} - C \left( \mathbf{n}^2 - \frac{1}{3} \mathbf{I} \right) + \frac{1}{3} D \mathbf{I} \quad (21)$ $B = 1 + \frac{3(1-c_2)}{2c_2} g_2(\theta, c_2) \cos 3\theta,$	Dafalias et al.(2004), Loukidis & Salgado (2009)
18 19 20 21 22 23 24	$c_2$	Parameter affecting the slope of the plastic potential	$C = 3 \sqrt{\frac{3}{2}} \frac{(1 - c_2)}{c_2} g_2(\theta, c_2)$ $g_2(\theta, c_2) = \frac{2c_2}{(1+c_2)-(1-c_2) \cos 3\theta} \quad (22)$ $D = A_o (\alpha_0^d - \alpha) : \mathbf{n} \quad (23)$	
25 26 27 28 29 30 31 32 33 34 35 36 37	$a$	Parameter controlling the initial fabric	$F_{11} = F_{22} = 0.5(1 - a), F_{33} = a,$ $F_{ij} = 0 \quad (\text{when } i \neq j) \quad (24)^{\#}$ $A = g_1(\theta, c_1) \mathbf{F} : \mathbf{n}, \quad (25)$ $A_c = - \left( \frac{1}{c_1} \right) A_e = \sqrt{\frac{3}{2}} \left( a - \frac{1}{3} \right)$ $A_{ratio} = 2 * \frac{A_e - A}{A_e - A_c} \quad (26)$	Dafalias et al.(2004)
38 39 40 41 42 43 44 45 46 47 48 49	$b_1, b_2, k$	Factors to scale the state parameter and plastic modulus for loading conditions other than triaxial compression	$f = \frac{1}{b_1 + b_2 (A_{ratio})^{b_3}} \quad (27)^*$ $b_3 = \frac{\log\left(\frac{1 - b_1}{b_2}\right)}{\log 2}$ $f_h = f^k \quad (28)^*$ $\psi_f = f \psi \quad (29)^*$	

\* Introduced for state parameter scaling

#  $F_{33}$  is assumed along the direction parallel to axis of particle deposition

Table 4 : Model parameters for UWA silica sand

Sanisand Model parameters	Symbol	UWA silica sand
Elastic properties	$G_0$	135
	$\nu$	0.14
Critical state properties	$e_o$	0.812
	$\lambda$	0.0189
	$\zeta$	0.7
	$M$	1.296
	$c_1$	0.7*
Yield surface	$m$	0.05
Parameter for Plastic potential gradient	$c_2$	0.78
Parameters influencing plastic modulus	$h_0$	7.5
	$c_h$	1.01
	$n_b$	2.0
Parameters influencing dilatancy	$A_0$	0.84
	$n_d$	3.4
Anisotropy and state parameter scaling parameters	$a$	0.27
	$b_1$	2.5
	$b_2$	-0.833*
	$k$	1.5

\*Assumed

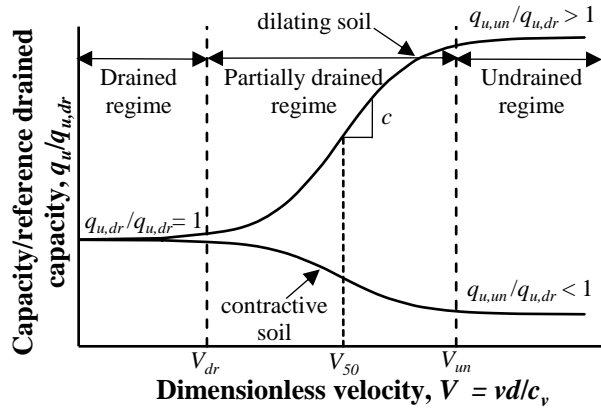


Fig. 1 : Backbone curve framework showing the variation in resistance across different consolidation regimes

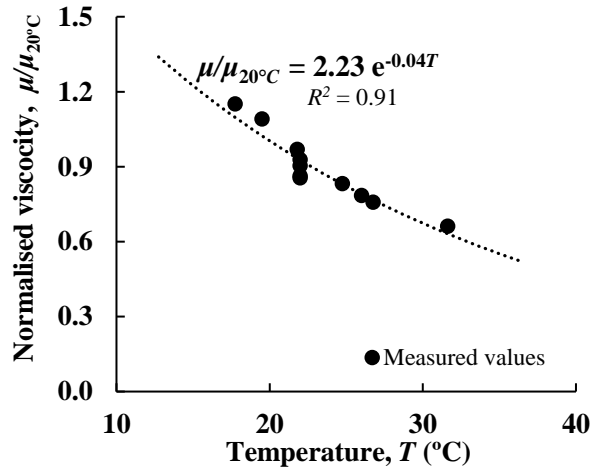


Fig. 2 : Variation of normalised viscosity with temperature for methocel F450

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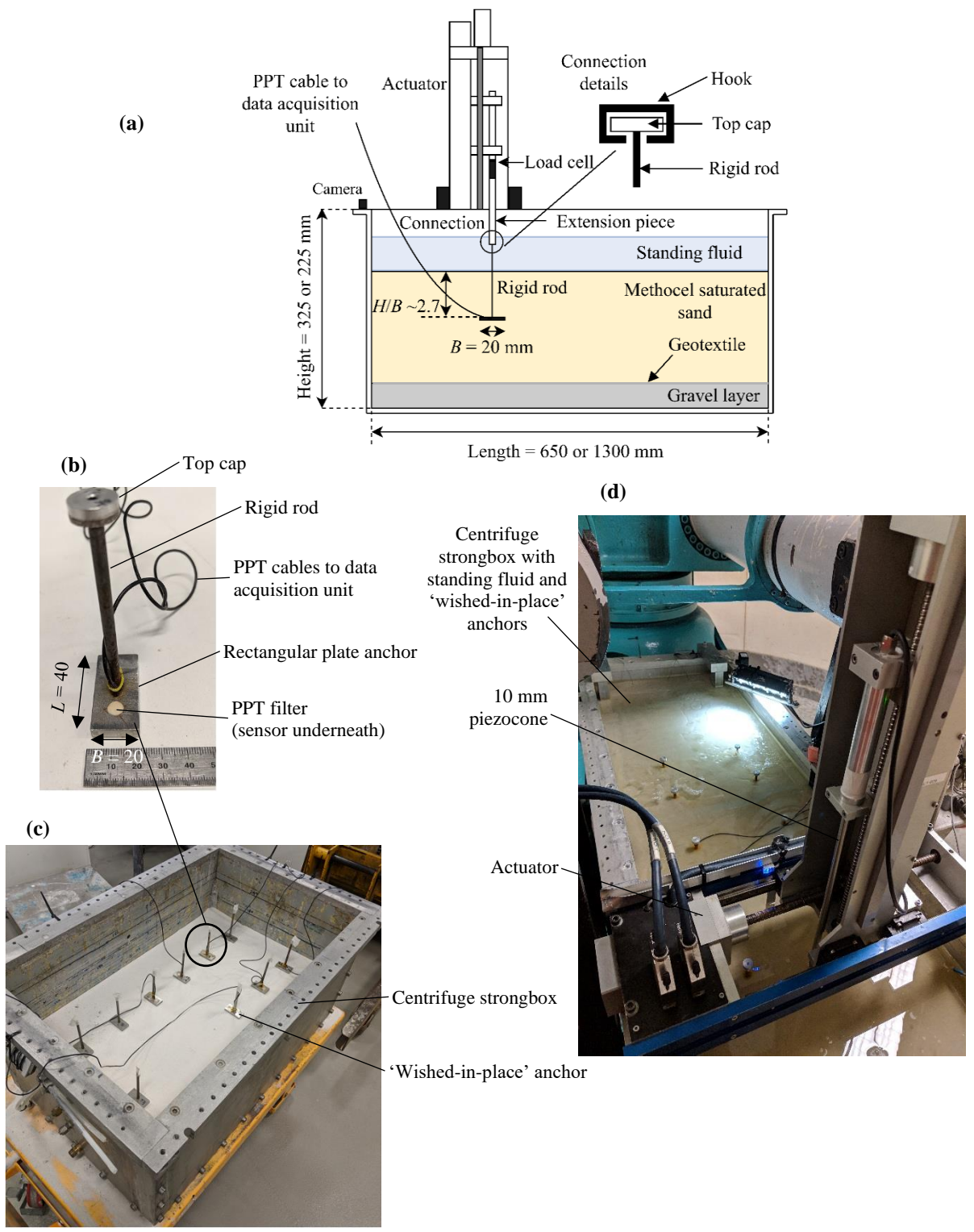


Fig. 3 : Centrifuge test details: (a) test setup (not to scale), (b) plate anchor model, (c) “wished in place” anchors during sample preparation (d) piezocone testing in saturated sand for sample characterisation

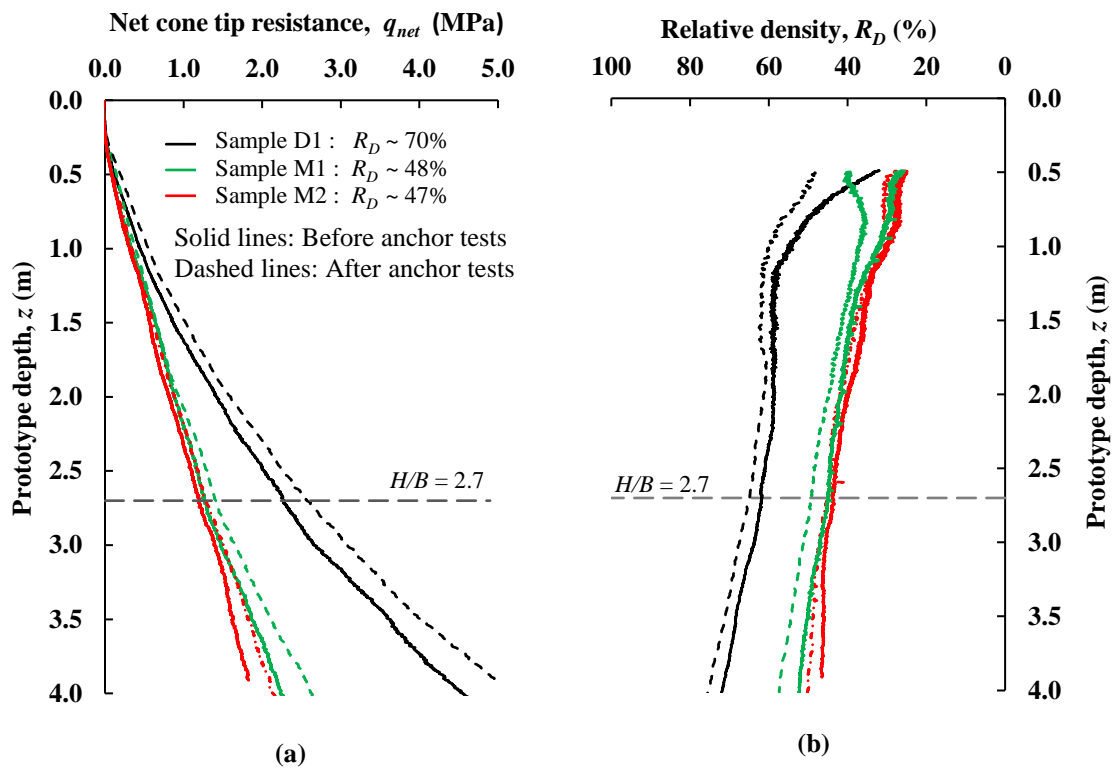


Fig. 4 : Profiles of: (a) net cone tip resistance and (b) relative density estimated using the empirical correlations in Roy et al. (2019)

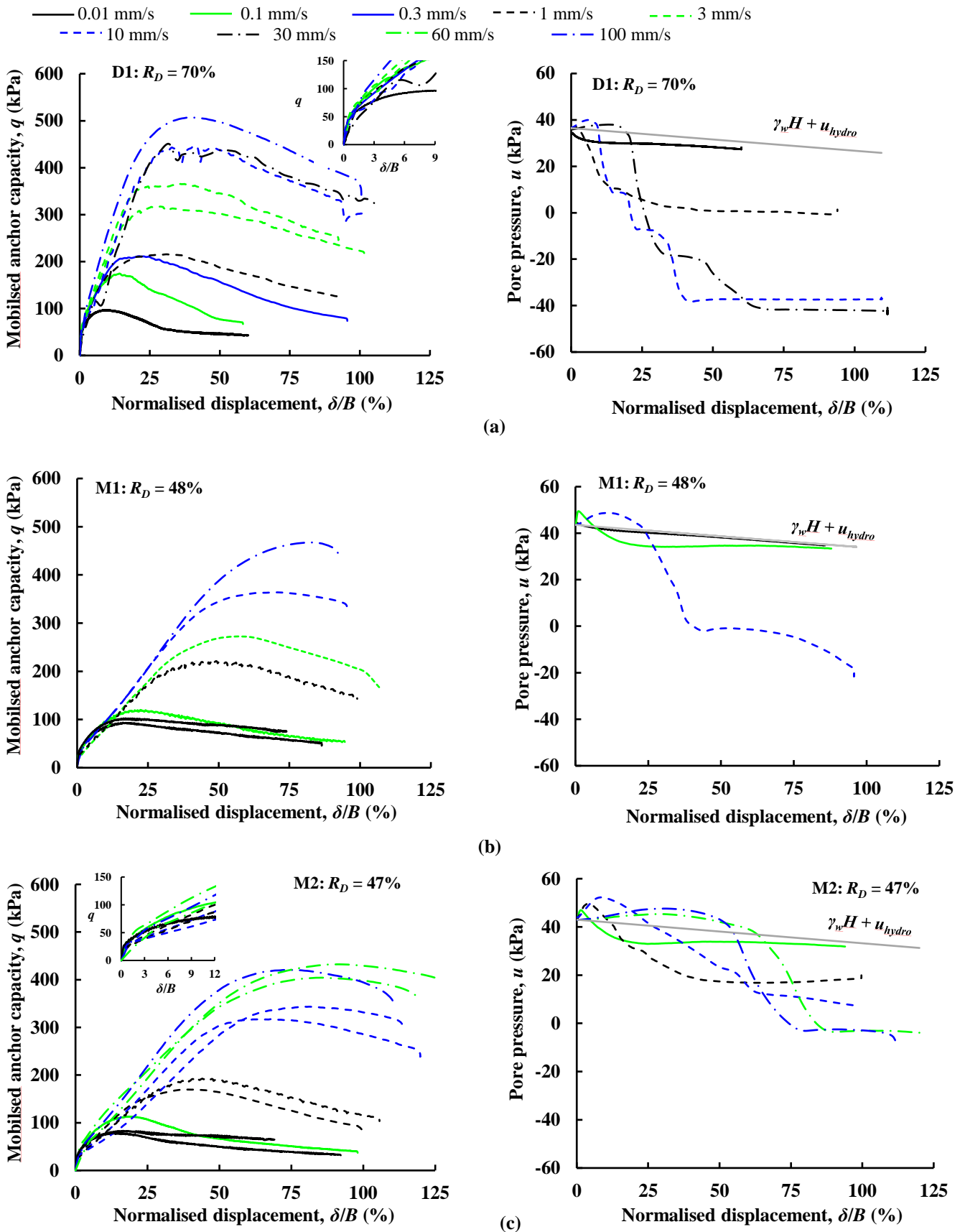
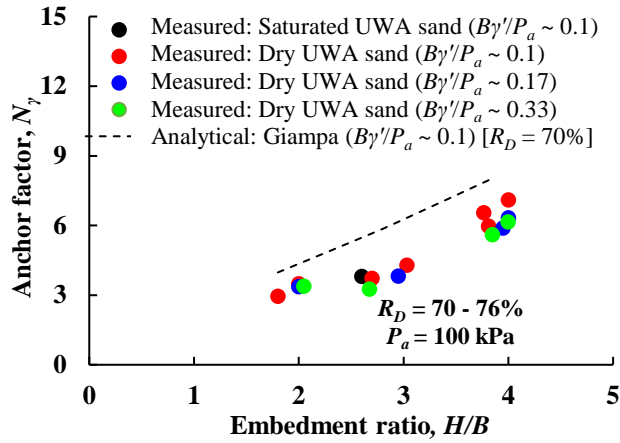
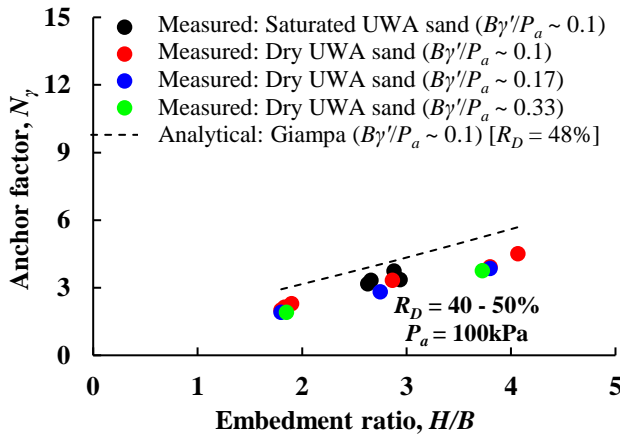


Fig. 5 : Mobilised capacity and pore pressure response at different uplift velocities in (a) dense sand of  $R_D = 70\%$  (b) medium dense sand of  $R_D = 48\%$  (c) medium dense sand of  $R_D = 47\%$

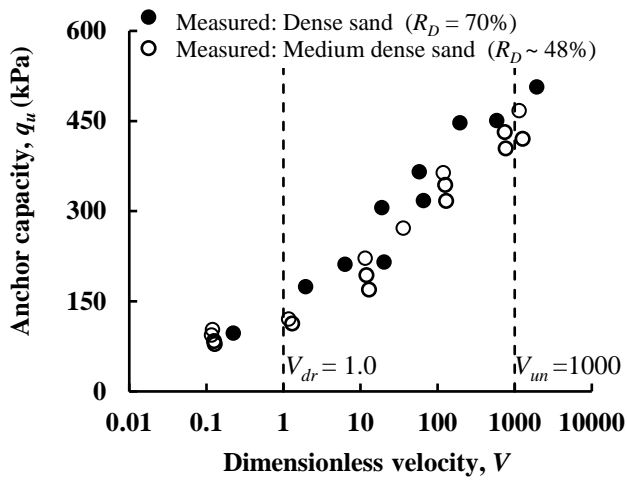


(a)

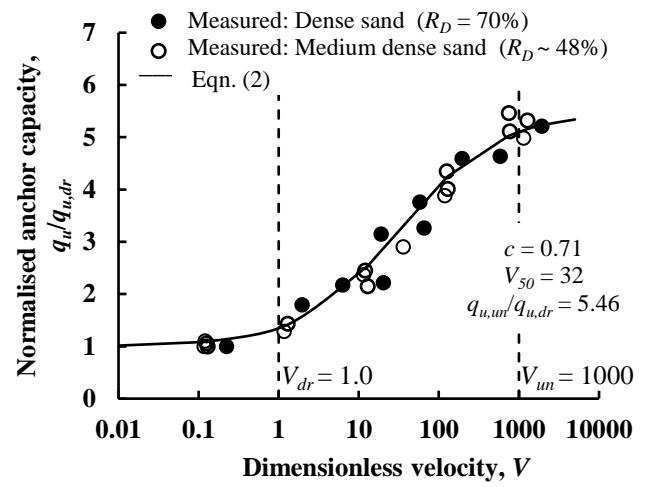


(b)

Fig. 6 : Comparison of measured  $N_\gamma$  for rectangular plates in dry and saturated drained conditions for (a) dense sand of  $R_D = 70\% - 76\%$  (b) medium dense sand of  $R_D = 40\% - 50\%$



(a)



(b)

Fig. 7 : Changing anchor capacity with dimensionless velocity: (a) mobilised anchor capacity ( $q_u$ ) and (b) normalised mobilised anchor capacity ( $q_u/q_{u,dr}$ )

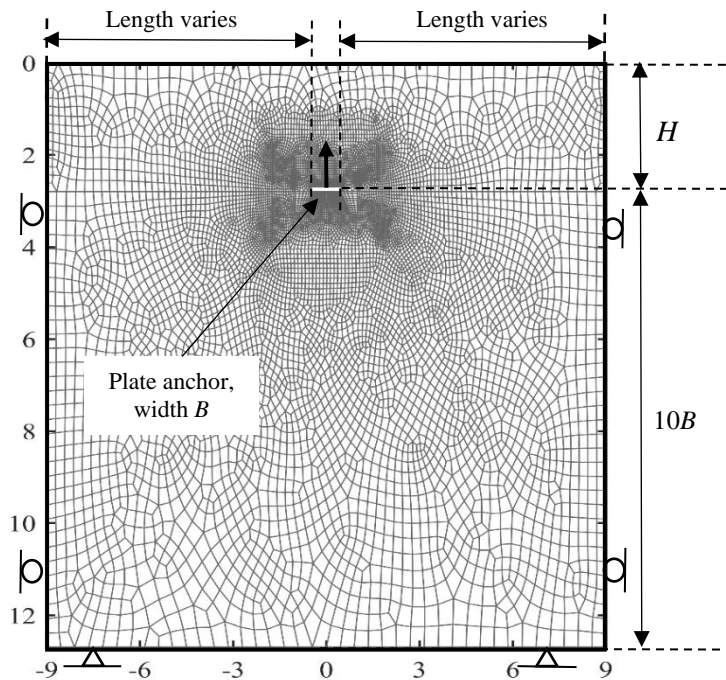


Fig. 8: Typical mesh details and boundary conditions for the anchor problem in Abaqus ( $H/B = 2.7$ )

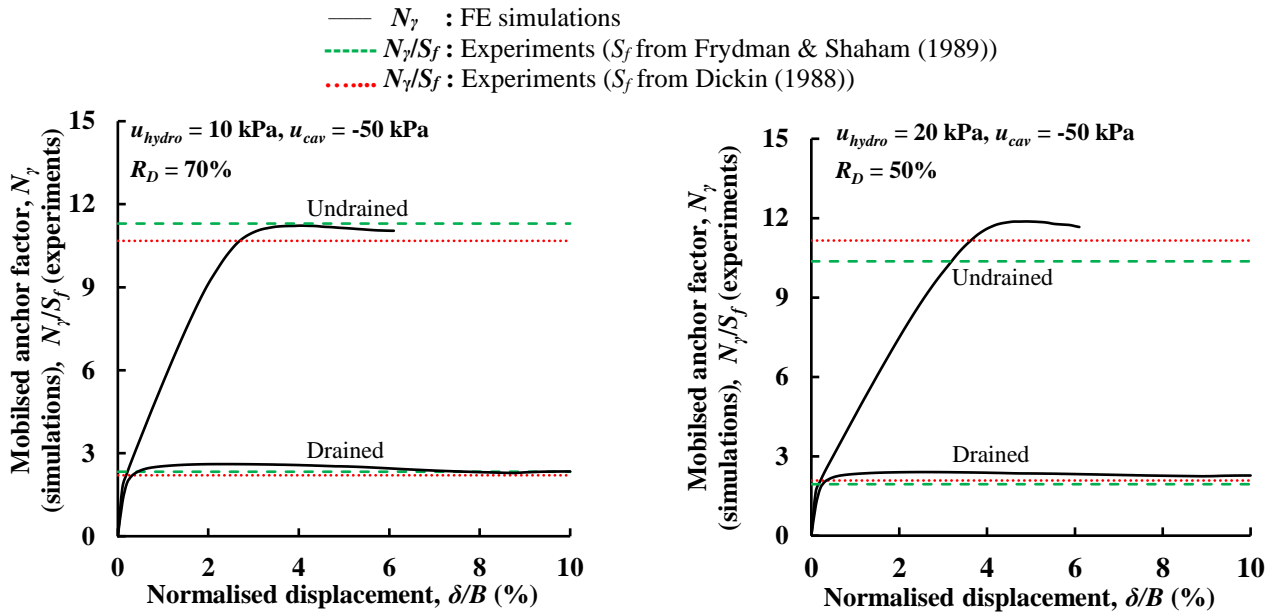


Fig. 9: Comparison of simulated and measured anchor factors (adjusted by shape factors) at  $H/B = 2.7$  in: (a) dense sand ( $R_D = 70\%$ ) (b) medium dense sand ( $R_D = 50\%$ )

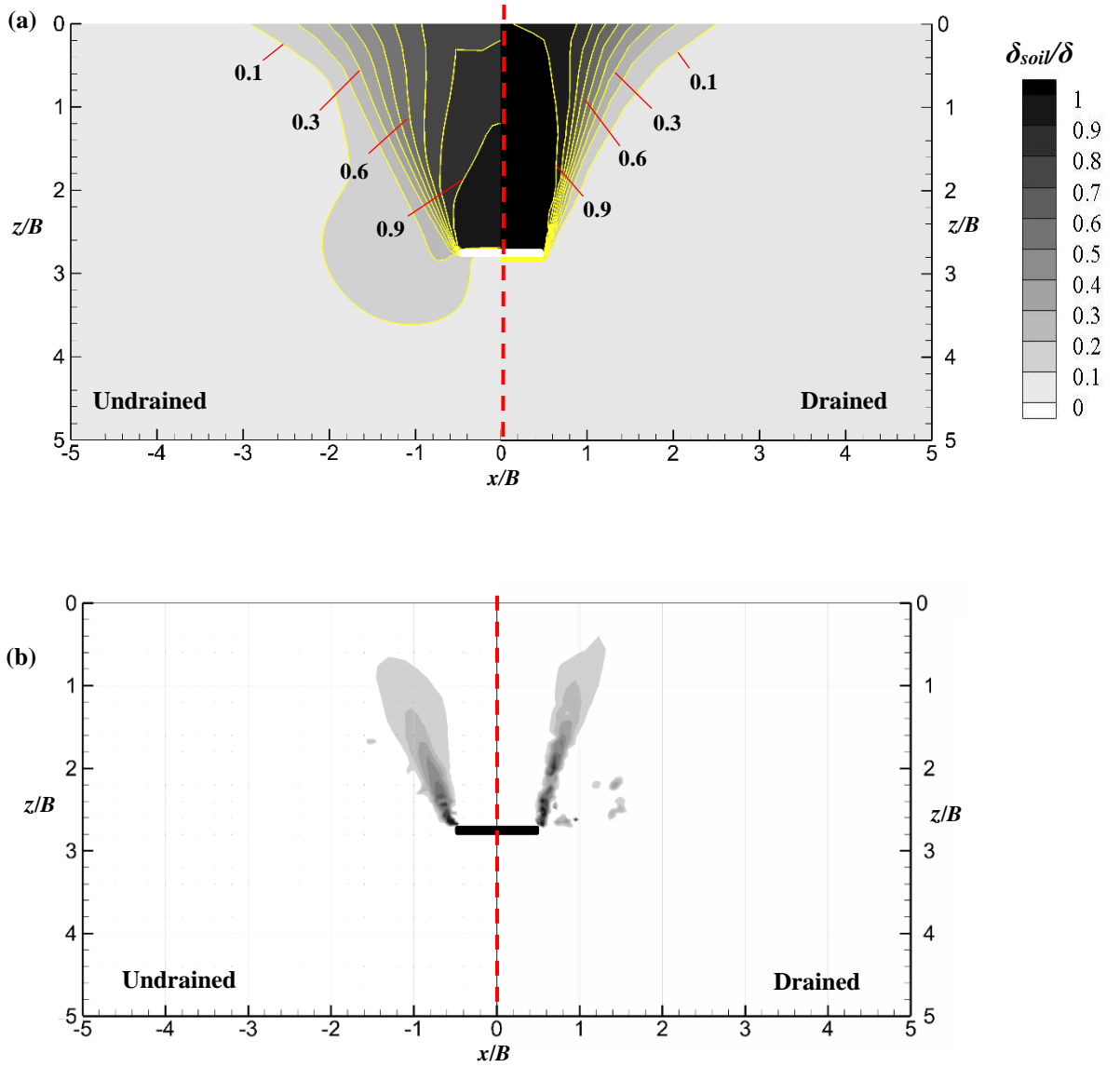
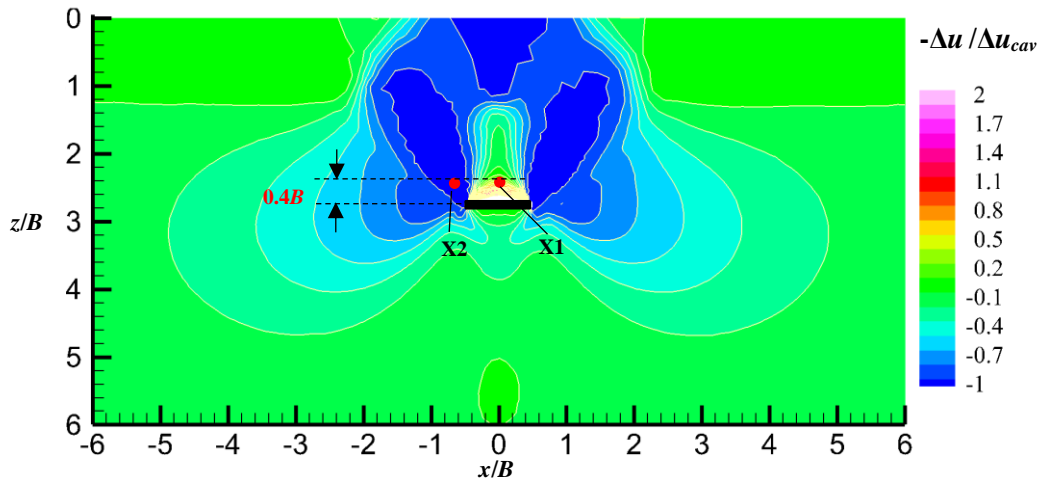
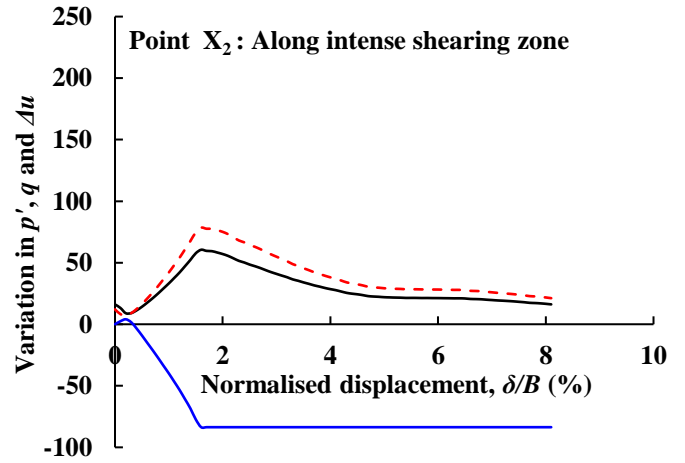
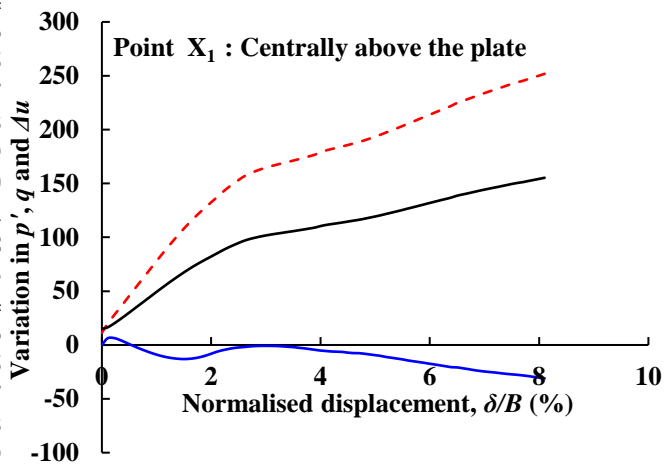
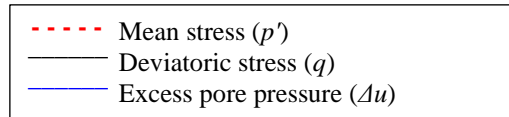


Fig. 10: Comparison of numerically determined: (a) failure mechanisms, (b) shear strain distributions at  $H/B = 2.7$  in dense sand ( $R_D = 70\%$ ) for undrained and drained cases



(a)



(b)

Fig. 11: (a) contours of excess pore pressure ( $\Delta u$ ) normalised by the maximum possible pore pressure change possible prior to cavitation ( $\Delta u_{cav}$ ) at  $H/B = 2.7$  in  $R_D = 70\%$  with  $u_{hydro} = 10$  kPa (b) variation of  $p'$ ,  $q$  and  $\Delta u$  at two different points in the domain

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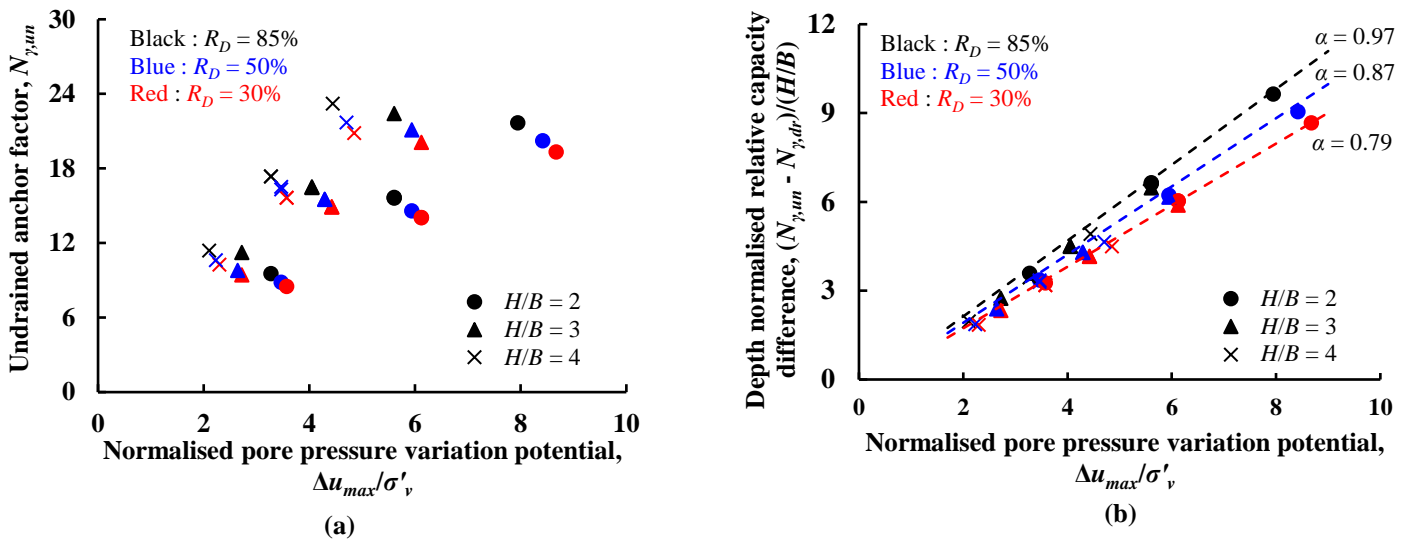


Fig. 12: Variation of: (a)  $N_{\gamma,un}$  and (b) depth normalised relative difference,  $\frac{N_{\gamma,un} - N_{\gamma,dr}}{H/B}$  from the FE simulations in dense and medium dense sands

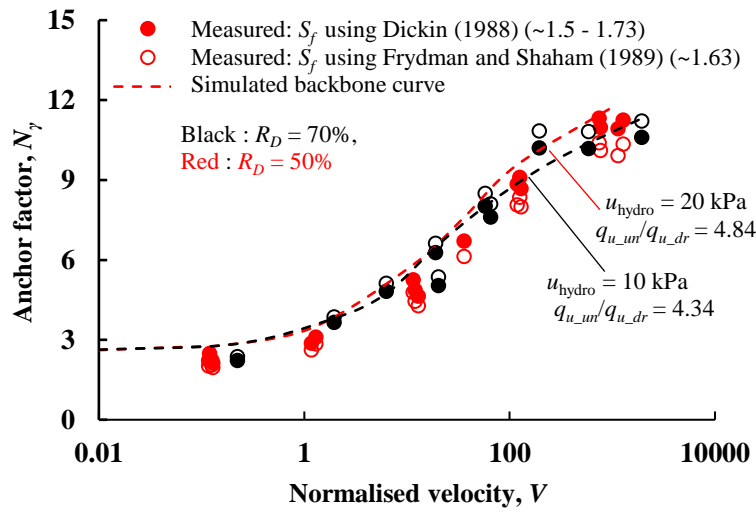


Fig. 13: Comparison of measured and predicted anchor factors with changing dimensionless velocity in dense and medium dense sand

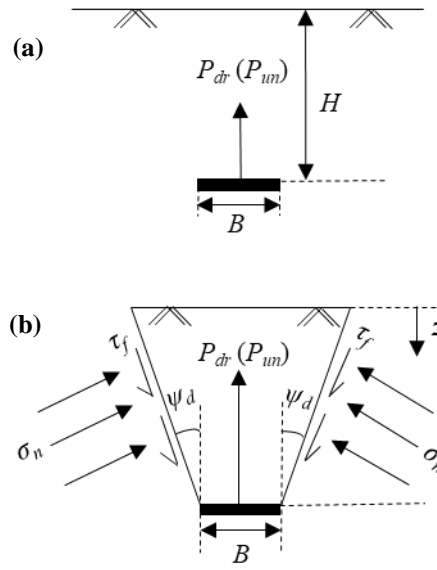


Fig. A1: Schematic representation of: (a) plate anchor uplift at failure (b) failure mechanism with inclined slip surface